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THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES.(U)

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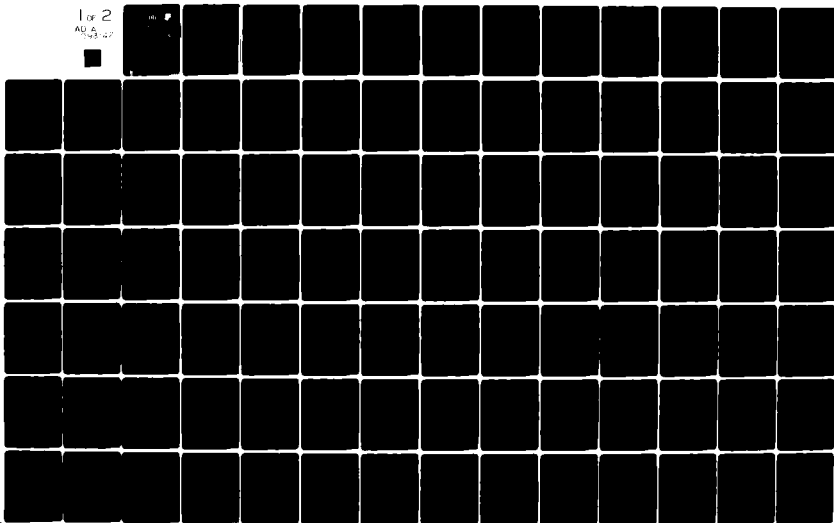
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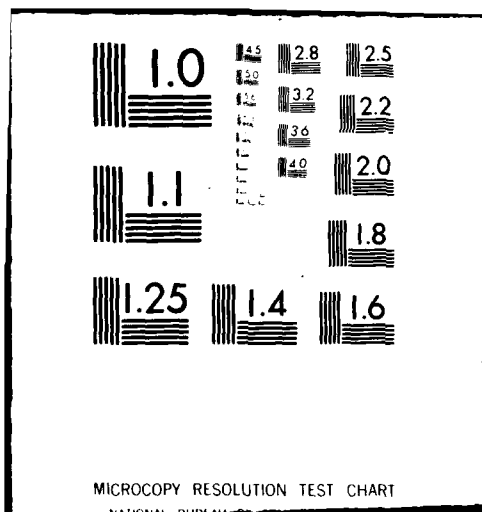
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Final Technical Report
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THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES

EMTEC Engineering Incorporated

C. Yeh

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This report has been reviewed by the RADC Public Affairs Office (PA) and is releasable to the National Technical Information Service (NTIS). At NTIS it will be releasable to the general public, including foreign nations.

Some referenced figures do not appear in this document. The references are to computer calculations too voluminous to publish. The resulting analysis of the computer data is adequately presented in this report. Therefore, the missing data is considered irrelevant to the conclusions presented herein. The missing data may be obtained by contacting RADC (ESO) Hanscom AFB MA 01731.

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→ treat problems dealing with nonlinear fibers or fibers with frozen-in statistically varying index profiles.

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EVALUATION

The effort as summarized in the report provides accurate theoretical analyses of the optical power transmission properties of a number of devices important to RADC efforts under TPO #4/D - Solid State Devices, Subthrust #3 - Electro-Optical Components. The devices include couplers, tapers, horns, and branches. The programs and techniques provided by the contractor permit in effect computer experiments to be done for a very large variety of design parameters. To do the actual experiments in the laboratory with this range of parameters would be enormously more expensive and time consuming. The report is then a crucial element in simplifying and accelerating the design process and in leading to final design specs in the shortest time.


LEONARD J. EYGES
Project Engineer

I. INTRODUCTION

This final report summarizes the work performed under Contract F19628-80-C-0053 which the Electronic Systems Division of the Air Force Systems Command granted to the EMtec Engineering, Los Angeles, California. The work was begun in January, 1980 and completed in August, 1980.

The principal thrusts of this R & D study in performing numerical analysis of multimode fiber components were two fold. Firstly, we wish to learn the limitation (and possible improvement) of our numerical scheme¹ and secondly, we wish to obtain numerical data for realistic multimode fiber structures.

Specifically, the following tasks were carried out:

- a) Study the effect of step index gradient and of tight beam confinement by an adaptive coordinate scheme.
- b) Study the effect of the presence of absorber at the edge of the mesh on the beam propagation characteristics of multimode fiber structures.
- c) Compute the coupling characteristics of tapered multimode fiber couplers and unequal size fiber couplers .
- d) Obtain data for reflection coefficients and beam waist changes for multimode fiber tapers, horns and branches.

In section II we shall present the implementation of the adaptive coordinates in our numerical solution of the scalar wave equation. Then, the scheme to include an absorber at the edge of the mesh will be described. Finally, an approximate approach to obtain the reflection coefficients for complex fiber structures will be shown. Detailed results of our study on the proposed tasks are given in Section III. Concluding remarks and recommendations for future work are included in Section IV.

II. ANALYTICAL APPROACH

The basic approach taken to find the solution of wave propagation along complex fiber structures is to solve the reduced scalar wave equation via the fast Fourier transform (FFT) technique². In this section we shall first indicate the conditions underwhich the exact vector wave equation may be simplified to yield the reduced scalar wave equation. Then we shall introduce the concept of adaptive coordinates³ and incorporate this concept in the solution of the reduced scalar wave equation via the fast Fourier transform technique.

A. Formulation of the Scalar Wave Approach. Starting with the vector wave equation for the electric field vector \underline{E} in the fiber structure,

$$\nabla \times \nabla \times \underline{E} - \omega^2 \mu_0 \epsilon \underline{E} = 0 \quad (1)$$

where ω is the frequency of the wave, μ_0 the permeability and $\epsilon = \epsilon(r)$, the inhomogeneous permittivity of the structure, and making use of the vector identity

$$\nabla \times \nabla \times \underline{E} = \nabla(\nabla \cdot \underline{E}) - \nabla^2 \underline{E} \quad (2)$$

and the relation

$$\nabla \cdot \underline{E} = -\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E}, \quad (3)$$

one has

$$\nabla^2 \underline{E} + \omega^2 \mu_0 \epsilon \underline{E} - \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) = 0 \quad (4)$$

Rewriting Eq. (4) gives

$$\nabla^2 \underline{E} + \omega^2 \mu_0 \epsilon_0 \left\{ \frac{\epsilon}{\epsilon_0} \underline{E} - \left[\frac{1}{\omega^2 \mu_0 \epsilon_0} \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) \right] \right\} = 0$$

The relative importance of the terms within the curly brackets can be determined from the following

$$\frac{\epsilon}{\epsilon_0} \underline{E} = \mathcal{O}\left(\frac{\epsilon}{\epsilon_0} \underline{E}\right) \quad (6)$$

$$\frac{1}{\omega^2 \mu \epsilon_0} \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E} \right) = \frac{1}{k_0^2} \mathcal{O}\left(\frac{\nabla \epsilon}{\epsilon} \cdot \nabla \underline{E}\right) = \mathcal{O}\left(\frac{\epsilon/\epsilon_0}{k_0 l} \underline{E}\right) \quad (7)$$

where the symbol \mathcal{O} means the "order of magnitude," and l is the smaller of the distance over which ϵ/ϵ_0 and \underline{E} change appreciably. For single-mode fiber structures, the values of ϵ/ϵ_0 and $k_0 \lambda$ are typically in the range

$$\epsilon/\epsilon_0 = \mathcal{O}(2) \quad (8)$$

$$k_0 l = \frac{2\pi}{\lambda} l = \mathcal{O}(10^2 \text{ or } 10^3), \quad l \approx \mathcal{O}(10\mu \text{ to } 100) \quad (9)$$

$$\lambda = \mathcal{O}(1\mu)$$

It follows that the second term within the curly brackets in Eq. (5) is several orders of magnitude smaller than the first term $\epsilon/\epsilon_0 \underline{E}$. It is therefore justifiable to neglect the second term and write Eq. (5) in the form

$$\nabla^2 \underline{E} + k_0^2 \frac{\epsilon}{\epsilon_0} \underline{E} = 0$$

The physical significance of replacing Eq. (5) by Eq. (10) is this. By discarding the term $\nabla \frac{1}{\epsilon} \nabla \epsilon \cdot \underline{E}$, we are neglecting any depolarization effects that may occur. This means that the wave retains the polarization it has at the source, which is evidenced by the fact that Eq. (10) can be reduced to a scalar equation by writing $\underline{E}(x)$ in the form

$$\underline{E}(\underline{x}) = \underline{e}_p u(\underline{x}) \quad (11)$$

where \underline{e}_p is a unit vector in the direction of the initial polarization of the wave.⁴ Substituting Eq. (11) in Eq. (10), we find that $u(\underline{x})$ satisfies the scalar wave equation,

$$\nabla^2 u + k_o^2 \frac{\epsilon}{\epsilon_o} u = 0 \quad (12)$$

This equation with the boundary condition on the initial surface, and the radiation condition at infinity, completely specifies $u(\underline{x})$, from which we can then obtain the electromagnetic field vectors \underline{E} and \underline{H} .

If we write u as the product of a factor $e^{ik_o z}$ that accounts for the rapid change in the phase of u along the direction of propagation and a complex amplitude $A(\underline{x}, z)$, a further simplification of the problem results

$$[2ik_o \frac{\partial}{\partial z} + \nabla_T^2 + k^2(n^2(\underline{x}, z) - n_o^2)] A(\underline{x}, z) = -\frac{\partial^2 A(\underline{x}, z)}{\partial z^2} \quad (13)$$

where ∇_T^2 is the transverse Laplacian $\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}$, and n_o is a given constant which represents the refractive index of some uniform medium. At laser wavelengths the complex amplitude $A(\underline{x})$ varies much more rapidly transverse to the direction of propagation than it does along the direction of propagation. This enables us to make the paraxial approximation wherein the term on the right side of Eq. (13) is neglected (in the Russian literature this is called the parabolic approximation). So, the complex amplitude now satisfies

$$\left[i2kn_o \frac{\partial}{\partial z} + \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + k^2(n^2(\underline{x}, z) - n_o^2) \right] A(\underline{x}, z) = 0 \quad (14)$$

For given initial data, i.e., values of the field at points on the initial surface, the propagation simulator must generate the corresponding field values at the terminal aperture such that Eq. (14) is satisfied. To do this we divide the medium into slabs defined by planes on which z is constant. In going from one slab to the next, we write $A(\underline{x}, z)$ in the form

$$A(\underline{x}, z) = e^{\Gamma(\underline{x}, z)} w(\underline{x}, z) \quad (15)$$

where $\Gamma(\underline{x}, z)$ is a phase function associated with the medium inhomogeneities

$$\Gamma(\underline{x}, z) = \frac{1k}{2} \int_{z_0}^z \left[n^2(\underline{x}, y, z') - n_0^2 \right] dz'. \quad (16)$$

The modified complex amplitude $w(\underline{x}, z)$ then satisfies the equation

$$\left[i2kn_0 \frac{\partial}{\partial z} + e^{-\Gamma} \nabla_T^2 e^{\Gamma} \right] w(\underline{x}, z) = 0 \quad (17)$$

with the initial condition

$$w(\underline{x}, y, 0) = u(\underline{x}, y, 0) \quad (18)$$

Physically, these equations approximate the propagation in the inhomogeneous medium by a two-step process at each z increment. First, we propagate the field $u(\underline{x})$ at $z - \Delta z/2$ to $z + \Delta z/2$, assuming that the intervening space is homogeneous. The effect of the inhomogeneities between $z - \Delta z/2$ and $z + \Delta z/2$ is then accounted for by multiplying this solution by the phase factor $\exp(\Gamma)$.

B. Adaptive Coordinates

To reduce the size of the mesh required to solve Eq. (17) numerically,

let us introduce an adaptive coordinate system defined by the transformation⁵

$$\zeta_1 = \frac{x/\rho_0}{N(z)} \quad (19)$$

$$\zeta_2 = \frac{y/\rho_0}{N(z)} \quad (20)$$

$$N(z) = \alpha^{-1/2} \left[\left(1 - \frac{z}{f}\right)^2 + \alpha^2 \left(\frac{z}{k\rho_0^2}\right)^2 \right]^{1/2} \quad (21)$$

$$\xi = \tan^{-1} \left[\frac{(1+\beta) \frac{z}{f}}{\beta^{1/2}} - 1 \right] \quad (22)$$

$$\beta^{1/2} = \alpha \frac{f}{k\rho_0^2} \quad (23)$$

where ρ_0 is a characteristic dimension of the beam at the initial surface (e.g., the e-folding radius of a gaussian beam), f is the distance to the focus, and α is a constant determined by the requirement that the solution be confined within the boundaries of the mesh at the focal plane. The choice $\alpha = 1$ yields a coordinate system that converges at a rate determined by the free-space diffraction of a gaussian beam having an e-folding radius ρ_0 .

When written in terms of the converging coordinate variables defined above, Eqs. (15) and (17) for the complex amplitude are replaced by the relations

$$w(x, y, z) = \hat{w}(\zeta, \xi) \exp(\tilde{\Gamma}) v(\zeta, \xi) \quad (24)$$

$$\hat{w}(\zeta, \xi) = \left(\alpha^{1/2} N(z) \right)^{-1} \exp \left[\frac{i}{2} (\zeta_1^2 + \zeta_2^2) \tan \xi \right] \quad (25)$$

$$\tilde{\Gamma} = \frac{ik}{2} \int_{z-\Delta z/2}^{z+\Delta z/2} dz' \left(n^2(x, y, z') - 1 \right) - \frac{i}{2} (\zeta_1^2 + \zeta_2^2) \Delta \xi \quad (26)$$

$$\left[\frac{\partial}{\partial \xi} - \frac{i}{2} \exp(-\tilde{\Gamma}) \left(\frac{\partial^2}{\partial \zeta_1^2} + \frac{\partial^2}{\partial \zeta_2^2} \right) \exp(\tilde{\Gamma}) \right] v = 0 \quad (27)$$

where $\Delta\xi$ is the increment in ξ in going from $z-\Delta z/2$ to $z+\Delta z/2$. The initial condition for v is

$$v(\underline{z}, \xi) = w(x, y, z) / \hat{w}(\underline{z}, \xi) \quad (28)$$

To solve Eq. (27) we utilize the fact that for sufficiently small values of $\Delta\xi$ (i.e., Δz) the effect of the exponential factors $\exp(\pm \tilde{\Gamma})$ in this equation is small. Hence, we solve the simpler equation obtained when these factors are equated to unity

$$\left[\frac{\partial}{\partial \xi} - \frac{i}{2} \left(\frac{\partial^2}{\partial \zeta_1^2} + \frac{\partial^2}{\partial \zeta_2^2} \right) \right] v = 0 \quad (29)$$

We use a fast Fourier transform technique to solve Eq. (29). The basis of this approach is the fact that the solution of Eq. (29) can be expressed in the form of a discrete Fourier series

$$v(\underline{z}, \xi) = \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} v_{mn}(\xi, t) \exp[i(p_m \zeta_1 + q_n \zeta_2)] \quad (30)$$

where the Fourier coefficients v_{mn} are determined from the initial data and Eq. (29) as follows. The initial values of v_{mn} are obtained by taking the discrete Fourier transform of the initial values of $v(\underline{z}, \xi_1)$ over a mesh of points $\zeta_1 = [l - (N/2)] \Delta\zeta$, $\zeta_2 = [j - (N/2)] \Delta\zeta$ ($l, j = 0, 1, \dots, N-1$)

$$v_{mn}(\xi_1) = \frac{(-1)^{m+n}}{N^2} \sum_{l=0}^{N-1} \sum_{j=0}^{N-1} v\left(\left(l - \frac{N}{2}\right) \Delta\zeta, \left(j - \frac{N}{2}\right) \Delta\zeta, \xi_1\right) \exp\left[-\frac{i2\pi}{N} (ml + nj)\right] \quad (31)$$

The dependence of V_{mn} is then determined by substituting Eq. (30) in Eq. (31), which yields

$$\frac{\partial V_{mn}}{\partial \xi} + \frac{i}{2} (p_m^2 + q_n^2) V_{mn} = 0 \quad (32)$$

from which it follows that

$$V_{mn}(\xi) = V_{mn}(\xi_i) \exp \left[- \frac{i(p_m^2 + q_n^2) \Delta \xi}{2} \right] \quad (33)$$

Finally, it can be shown that in order for the discrete Fourier series representation of v given in Eq. (30) to be real when v is real, the coefficients p_m and q_n must have the form

$$p_m = \frac{2\pi}{N\Delta \zeta} \left(m - \frac{N}{2} \right) \quad (34)$$

$$q_n = \frac{2\pi}{N\Delta \zeta} \left(n - \frac{N}{2} \right) \quad (35)$$

Hence, for discrete points $\zeta_1 = (\ell - N/2)\Delta \zeta$, $\zeta_2 = (j - N/2)\Delta \zeta$ ($\ell, j = 0, 1, \dots, N-1$)

$$\begin{aligned} & v \left(\left(\ell - \frac{N}{2} \right) \Delta \zeta, \left(j - \frac{N}{2} \right) \Delta \zeta, \xi \right) \\ &= (-1)^{\ell+j} \sum_{m=0}^{N-1} \sum_{n=0}^{N-1} (-1)^{m+n} V_{mn}(\xi_i) \\ & \exp \left[- i \beta \left(\left(\frac{m - \frac{N}{2}}{N} \right)^2 + \left(\frac{n - \frac{N}{2}}{N} \right)^2 \right) + i \frac{2\pi}{N} (\ell m + j n) \right] \quad (36) \end{aligned}$$

where $\hat{\epsilon} = 2\pi^2 \Delta\xi / (\Delta\xi)^2$. Note that v is simply $(-1)^{l+j}$ times the discrete Fourier transform of $(-1)^{m+n} v_{mn}(\xi)$.

The effect of the medium and the factor $\exp[(-1/2(\xi_1^2 + \xi_2^2)\Delta\xi)]$ introduced by the coordinate transformation is taken into account at each ξ step in the calculation by multiplying the value of v obtained in the previous step by the quantity $\exp(\tilde{\Gamma})$ defined in Eq. (26), i.e., the initial value inserted in Eq. (31) is $\exp(\tilde{\Gamma})$ times the value of v determined from the previous steps.

Using this adaptive coordinate algorithm we have been successful in our treatment of various realistic multimode fiber structures. Results are summarized in Section III.

C. Implementation of Lossy Outer Boundary.

It is believed that the field touching the outer boundary of the cladding region of a fiber structure will be attenuated due to radiation or absorption. To accomodate this situation in order to further improve our computer simulation, we have incorporated the presence of a lossy dielectric layer outside the cladding region in our computer program. An example of the index profile of a fiber is shown in Fig. 1:

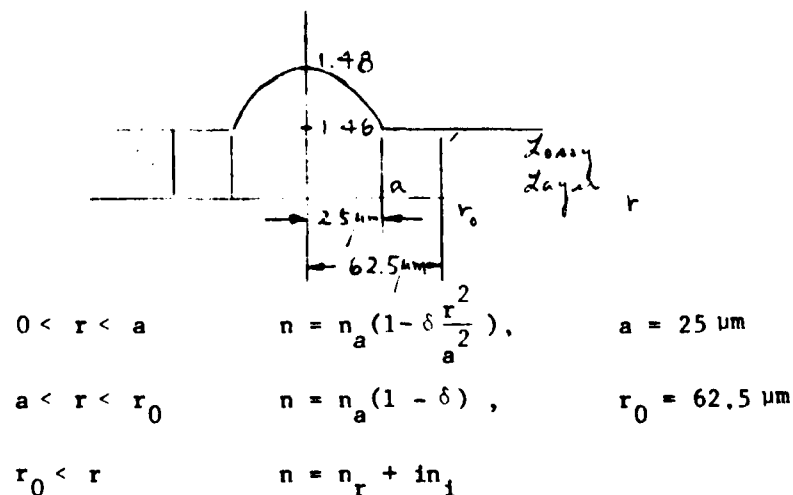


Figure 1: A Typical Index Profile with Lossy Outer Layer

Typical intensity patterns of beams propagating in a fiber with the index profile given by Fig. 1 are shown in Fig. 2. It was found that the propagation characteristics of the guided beams are not significantly affected by the presence of a lossy outer layer, except when the spot size of the beam is larger than the core diameter, as expected.

D. A Heuristic Approach in Obtaining the Reflection Coefficient.

One of the a'prior assumption in the development of the FFT scalar wave approach is that only paraxial rays are allowed and no reflection is permitted. This assumption enables us to develop an algorithm, thereby, we may obtain the propagating field by a forward stepping process as described earlier. As the field evolves from one z plane to the next $z + \Delta z$ plane, the averaged value of the refractive index as seen by the field may be different as illustrated in Fig. 3:

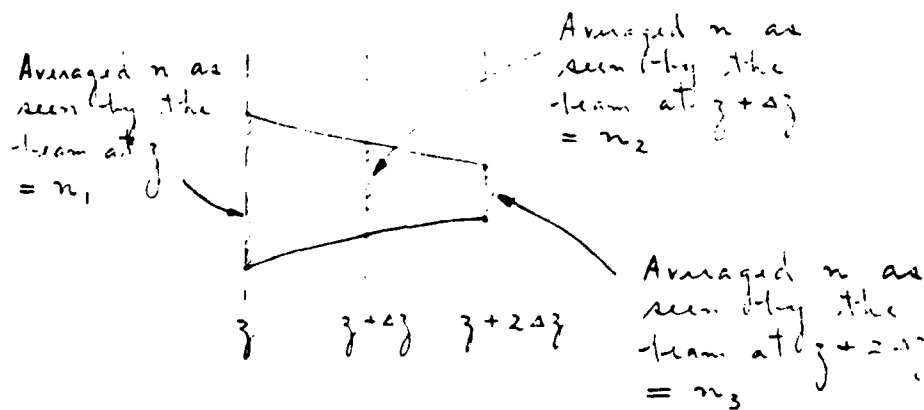
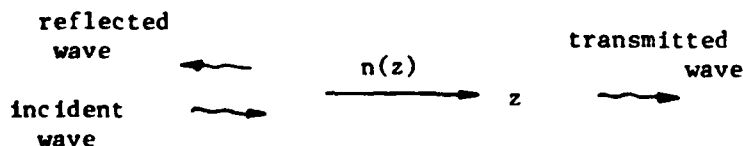


Figure 3: Illustration of the averaged n as seen by the beam.

In affect one may postulate that the wave is experiencing reflection in a medium with longitudinally slowly varying refractive index as shown below:



$n(z)$ is given by Fig. 3.

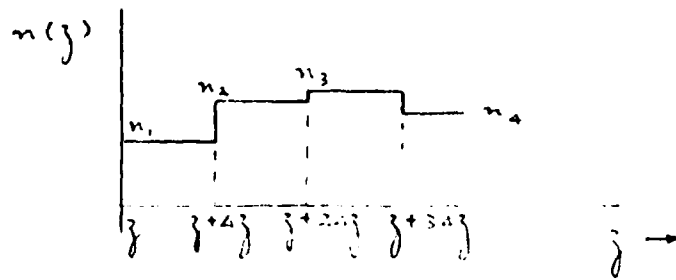


Figure 4: Equivalent Index Profile

The reflection coefficient for a plane wave propagating in this longitudinally non-uniform medium may be obtained according to a formula derived for the case of plane wave propagation in stratified layered medium:⁶

$$R(z) = -\exp[-is(z)] \int_z^\infty \gamma(z) \exp[is(z)] dz$$

$$s(z) = 2 \int^z \beta(z) dz \quad \beta(z) = k_0 n(z)$$

$$\gamma(z) = \frac{d\beta}{dz} / 2\beta$$

This is the heuristic approach that we shall use to calculate the reflection coefficient for waves in our multimode fiber structures.

III. Results

The algorithms detailed above have been implemented in our computer programs. Results for the proposed tasks are given in the following:

(a) Effects of Step Index Gradient on the Propagation Characteristics.

The purpose of this study is to learn the effects of step index gradient on the propagation characteristics of waves in a multimode fiber guide. Let us introduce the following index profile:

$$n(r) = n_0 - \delta \left(\frac{r}{a}\right)^{2m} \quad (0 < r < a)$$

$$n(r) = n_0 - \delta \quad (a < r)$$

where n_0 , m and δ are given constants and a is the core radius of the fiber. For a typical parabolic index profile fiber, one has

$$n_0 = 1.48, \quad \delta = 0.02, \quad a = 25\mu\text{m}, \quad m = 1.$$

The constant δ must necessarily be small so that the depolarization effect may be ignored and the scalar wave approach may be justified. By varying m , the steepness of the index gradient may be varied as shown in Fig. 5. It should be kept in mind that even when the FFT technique is capable of handling steep index variations, the slope of the index profile must still be gentle enough so that the gradient term in the exact wave equation (Eq. (5)) may be ignored. We have carried out propagation calculation for the following specific cases: $n_0 = 1.48$, $\delta = 0.02$, $a = 25\mu\text{m}$, $m = 1, 4, 6, 10$. Higher m values means steeper index gradient. As shown in Figs. 6, no computational difficulties were encountered for even the steepest case ($m = 10$) in which the index changes from $n = 1.48$ to $n = 1.46$ in $3\mu\text{m}$ distance for $50\mu\text{m}$ core diameter fiber. However, one should be aware that we are pushing the limit of validity for the scalar wave approach.

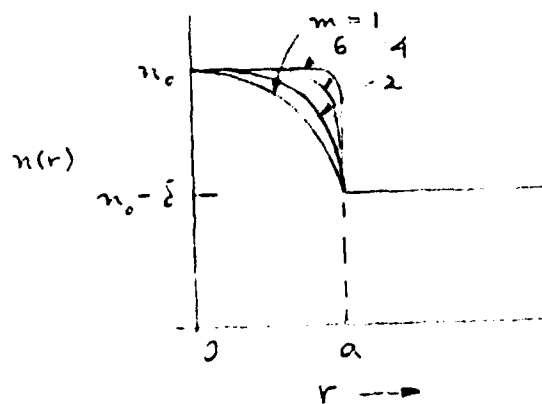


Fig. 5. Plot of $n(r) = n_0 - \delta (r/a)^{2m}$.

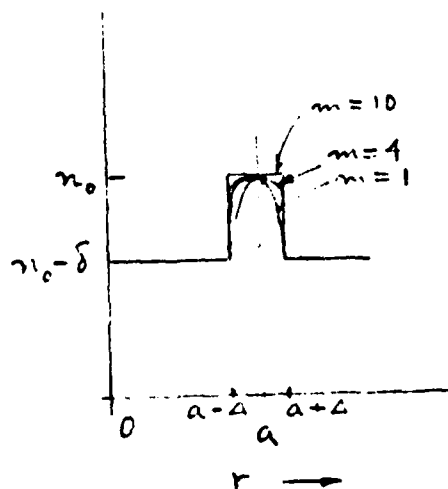


Fig 7 Plot of $n = n_0 - \delta ((r-a)/\Delta)^{2m}$.

From the results, it is of interest to note that as m increases from 1, i.e., as the index profile deviates from the parabolic profile, the beam profiles no longer remain to be of gaussian shapes, but take on ring-type structures. This implies that the phase front of the multimode beam is no longer a monotonic function of the radial distance but has become an oscillatory one.

We may conclude from these calculations that our program is capable of handling problems with steep index gradient. The index transition may occur in a distance as small as 4λ where λ is the free-space wavelength. The limiting factor apparently is the justification for the elimination of the depolarization term in Eq. (5).

(b) Beam Propagation in a Ring Fiber

A typical single-mode fiber has a core diameter of the order of $10\mu\text{m}$. Consequently it is more difficult to handle than a multi-mode fiber. An idea to enlarge the single-mode fiber has recently been put forth by Dr. L. Eyges of RADC. He suggested that perhaps a ring-type structure may support a single mode and yet possesses larger dimension than the usual solid-core fiber. This task was undertaken to investigate this possibility. Let us postulate that the index profile of a ring fiber takes the following form

$$n = n_0 - \delta \left(\frac{r-a}{\Delta}\right)^{2m} \quad \text{for } a-\Delta < r < a+\Delta$$

$$n = n_0 - \delta \quad \text{for } a-\Delta > r, r > a+\Delta$$

where n_0 , δ , a , Δ , and m are given constants. By increasing the m value, one may adjust the steepness of the index gradient as shown in Fig. 7. Two types of initial beam shapes will be studied: (1) a solid centered gaussian beam and (2) a hollow-centered donut beam. We wish to learn how well the ring fiber will confine these two types of beams. The field expression for a solid gaussian beam takes the form

$$u = e^{-\frac{\alpha}{2} \left(\frac{r}{a}\right)^2}$$

while the hollow-centered donut beam takes the form

$$u = e^{-\frac{\alpha}{2} \left[\left(\frac{r}{a} \right)^2 - 1 \right]}$$

where α and a are given constants. Results of our computation are shown in Figs. 8 and 9. By following the evolution of the beam intensities, one may determine how well the ring fiber is guiding the beam. It can be seen from these figures that the solid beams appear to be better confined than the hollow beams, although the spreading of the solid beam energy is quite noticeable. It also appears that simple insertion of beam energy in the high index region of the ring fiber does not insure good guidance of the beam energy. One may conclude from this preliminary study that neither solid Gaussian beams nor hollow-centered Gaussian beams correspond to the mode energy distribution of a single-mode in a ring fiber. One should first perform the classical modal analysis to obtain the mode pattern of the single mode and then use this mode pattern as the initial beam pattern for propagation down the ring fiber. It is believed that the use of ring-index fiber as large core single-mode fiber definitely possesses merit and should be studied further. What we have demonstrated with our present study is that our program is capable of handling this type of fibers.

(c) Fiber Couplers

One of the simplest type of light couplers is the fiber coupler. By placing two or more fibers in close proximity of each other light energy may transform from one to the other through the coupling effect. This coupling process is rather involved. The well-known coupled mode theory may be adequate for simple, single-mode structures such as slabs

with reasonable separations. But, when multi-mode complex structures such as the fiber couplers are involved, the coupled mode theory becomes grossly inadequate.⁷* On the other hand, our FFT-scalar wave approach is uniquely qualified to deal with this fiber coupler problem. This is because this technique provides the evolution of beam field as it propagates down a complex multimode inhomogeneous fiber structure. Four types of fiber couplers have been studied:

Case 1 Coupling between two equal parabolic index fibers.

Two graded-index fibers are fused together longitudinally with separation d between their centers. The index profile for each fiber is given by

$$n(r_{1,2}) = n_0 - \delta \left(\frac{r_{1,2}}{a} \right)^2$$

where n_0 , δ , and a are given constants, and 1 or 2 refers, respectively, to #1 or #2 fiber. Typical values for a Corning or ITT graded index fibers are used:

$$n_0 = 1.48$$

$$\delta = 0.02$$

$$a = 25\mu\text{m}$$

*Recent advances by L. Eyges and P. Gianino of RADC using the extended boundary condition technique have shown that single mode couplers involving arbitrarily shaped uniform core guides can be successfully and accurately treated.

Various separation d were used. A gaussian beam represented by

$$u(x,y) = u_0 \exp \left\{ \left[- \left(x + \frac{d}{2} \right)^2 - y^2 \right] / w^2 \right\}$$

where $u(x,y)$ is the scalar wave function of the beam, and u_0 , w are given constants, is incident on one of the fibers. Results have been obtained for

$$w = 2.5\mu m, 5\mu m, 10\mu m$$

$$d = 8\mu m, 12\mu m, 16\mu m, 20\mu m.$$

The evolution of the beam along this structure is shown in Figs. 10-11. Displayed in Figs. 12 is the %-power in one fiber as a function of longitudinal distance for various separations and initial beam sizes. For the situations considered above, many modes are excited. The coupling process is very involved as displayed in Fig. 10-11. It still appears that back and forth power exchange among the guides prevails. The complexity of the power exchange phenomenon for the multimode coupler re-emphasizes the importance of obtaining design data through analysis before the actual construction of fiber coupler.

Case 2 Coupling between two equal step-index fibers.

This coupler is identical to the previous one except step-index fibers were used. We shall approximate the index profile of a step index fiber by the following expression:

$a(0, -4), (0, 4)$
 Separations $b(0, -6), (0, 6)$
 $c(0, -8), (0, 8)$

Two Fiber Core
 Parabolic Index Profile
 Initial Beam Spot Size = $10\mu m$

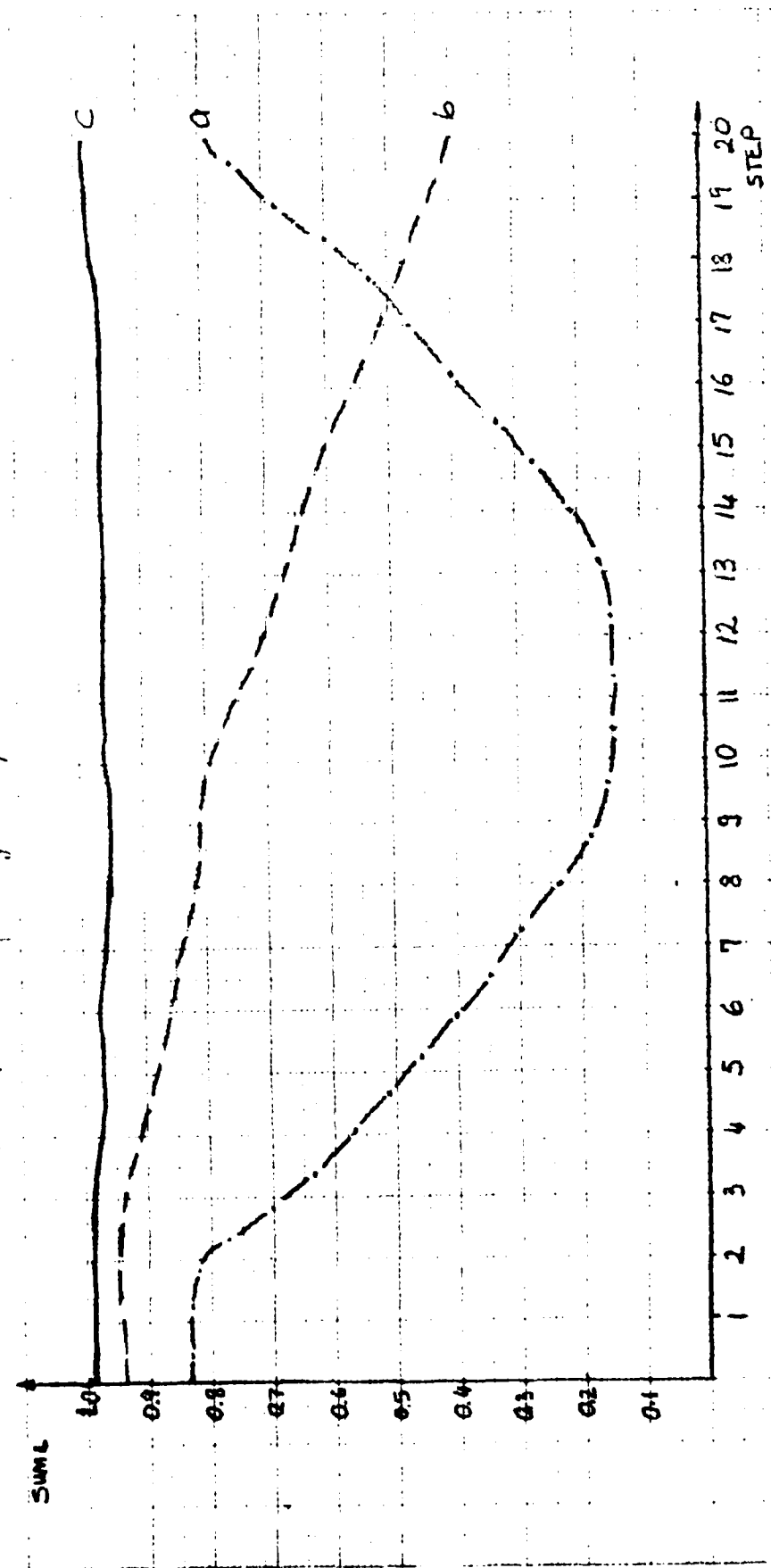


FIG 12 (a)

Pressure - 1400
 Beam spot size = 20 μ m

Two Fiber Coupler
 Parabolic Index Profile

a(0, -4), (0, 4)
 b(0, -6), (0, 6)
 c(0, -8), (0, 8)
 d(0, -11), (0, 11)

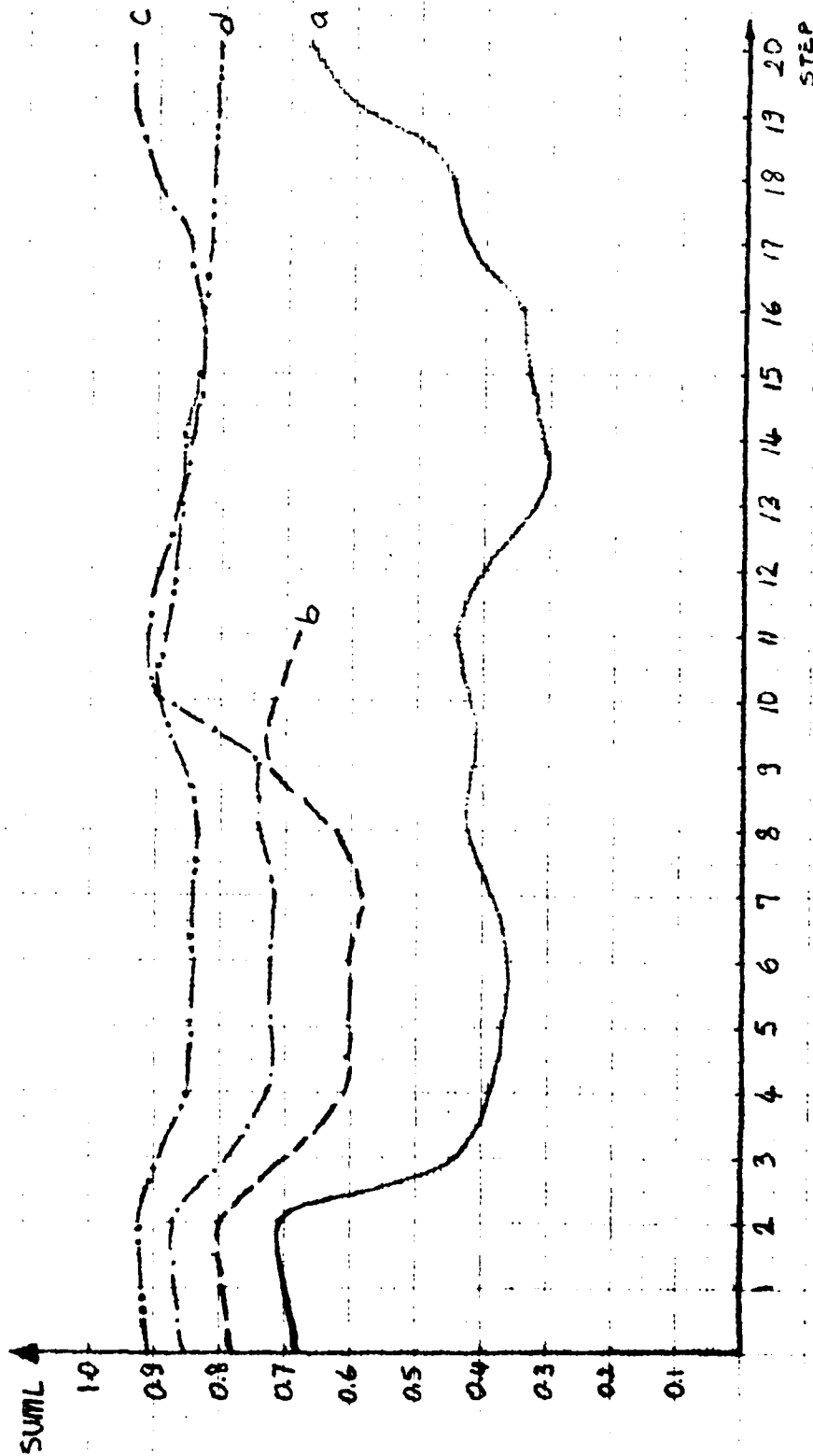


FIG 12 (b)

$$n(r_{1,2}) = n_0 - \delta \left(\frac{r_{1,2}}{a} \right)^{2m}$$

with $m = 4$. Again typical values for a Corning or ITT graded index fiber are used; i.e., $n_0 = 1.48$, $\delta = 0.02$, $a = 25\mu\text{m}$. Results have been obtained for

$$w = 2.5\mu\text{m}, 5\mu\text{m}$$

$$d = 8\mu\text{m}, 12\mu\text{m}, 16\mu\text{m}$$

where d is the separation distance and w is the beam waist radius. Specific results are given in Figures 13. Displayed in Figs. 14 is the % power in one fiber as a function of longitudinal distance for various separations.

Case 3 Coupling between two unequal fibers.

It is of interest to learn, when two unequal size fibers are placed side by side, whether transfer of power would occur for the multimode case. The following fibers were used:

$$n(r_{1,2}) = n_0 - \delta \left(\frac{r_{1,2}}{a_{1,2}} \right)^{2m}$$

$$n_0 = 1.48 \quad m = 1, 4$$

$$\delta = 0.02$$

$$a_1 = 25\mu\text{m}$$

$$a_2 = 12.5\mu\text{m}$$

$$\text{Separation distance, } d = 8\mu\text{m}, 12\mu\text{m}$$

$$\text{Initial Beam Radius, } w = 5\mu\text{m}, 10\mu\text{m}$$

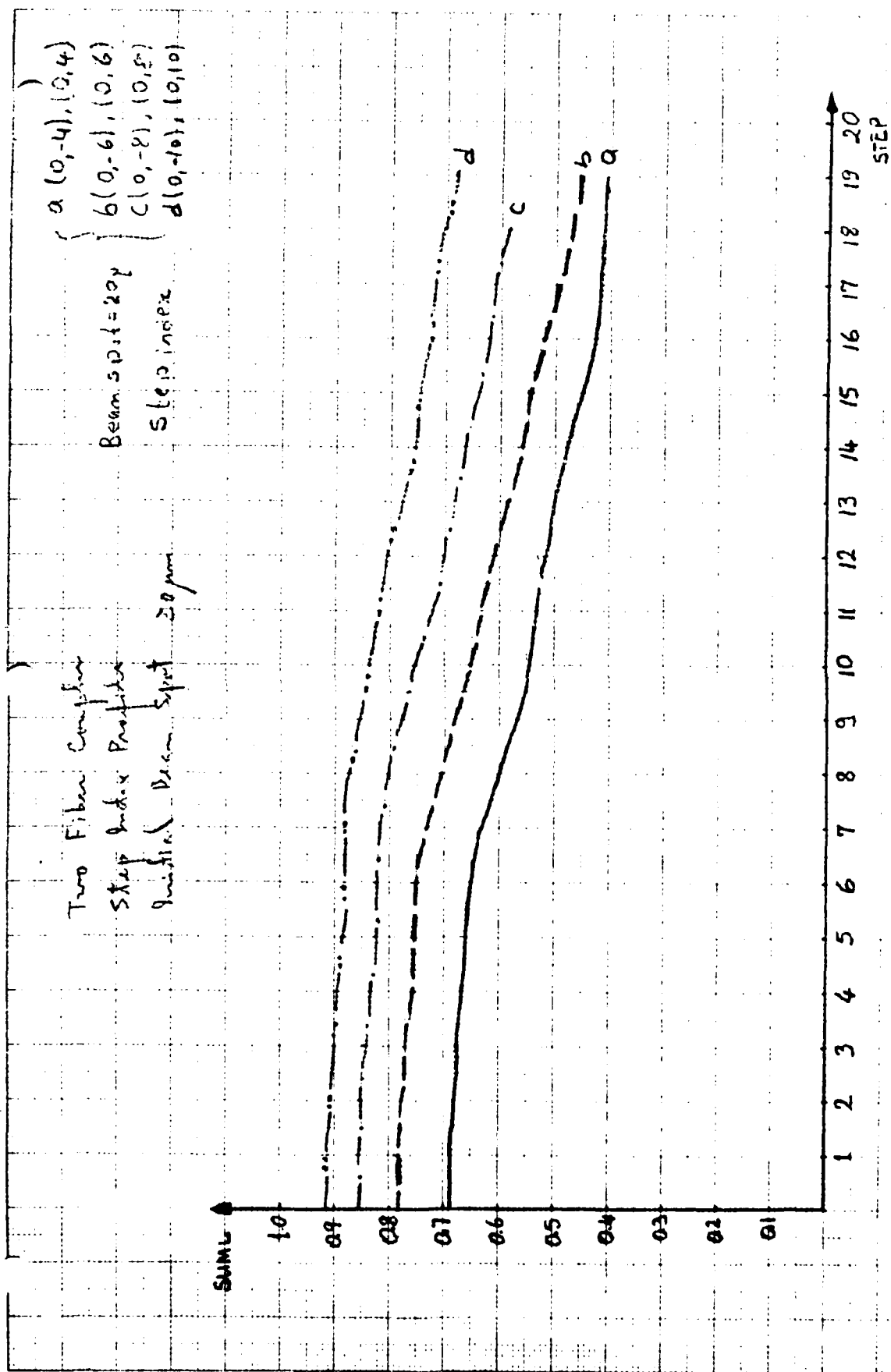


FIG 14

For the above chosen parameters, one may observe from Figs. 15-16 that only nominal coupling occurs between two unequal size fibers. In other words these structures are not efficient couplers.

Case 4 Coupling between more than two fibers.

As illustrations, we have considered two types of structures: Three equal size fibers located equal-distance from each other and three unequal size fibers arranged in a triangle shape. One fiber is initially illuminated, we wish to learn the power exchange characteristics in these two couplers. The following parameters were used:

$$n(r_{1,2,3}) = n_0 - \delta \left(\frac{r_{1,2,3}}{a_{1,2,3}} \right)^2$$

$$n_0 = 1.48$$

$$\delta = 0.02$$

$$a_1 = 25\mu\text{m}$$

$$a_2 = 25\mu\text{m}, 20\mu\text{m}$$

$$a_3 = 25\mu\text{m}, 15\mu\text{m}$$

$$\text{Separation distances, } d_1 = d_2 = d_3 = 12\mu\text{m}$$

$$\text{Initial beam radius, } w = 5\mu\text{m}$$

It can be seen from Fig.17 that when one of the fiber of 3 identical fibers, separated an equal distance from each other, is illuminated with a gaussian beam, power exchange takes place between two unilluminated fibers with the one illuminated fiber in a synchronous manner. In other words, each of the originally unilluminated fibers is receiving identical amount of power transfer from the illuminated one, as expected. On the other hand, the power exchange phenomenon for 3 non-identical fibers coupler is much more complicated as seen from Fig.18

Larger power transfer tends to take place among fibers with similar diameters.

Fiber coupler is one of the crucial components in any fiber optics system. As such, one must understand the detailed wave interaction phenomenon in this type of coupler so that correct design can be made. We have demonstrated the capability of our technique in dealing with this type of fiber structures. Systematic studies of different varieties of fiber couplers may now be undertaken.

(d) Reflection Coefficient Calculations.

We have implemented the heuristic approach discussed earlier in the computer program to yield reflection coefficients for waves propagating in the various multimode fiber structures. Although it is difficult to justify the accuracy of the absolute values for the reflection coefficients obtained according to this algorithm, nevertheless, we feel that their relative values for different structures can be believed. This is because our heuristic approach took into consideration the fundamental characteristic of wave reflection: i.e., reflection occurs when discontinuity of the propagation medium or structure is experienced by the wave. The larger is the discontinuity, the larger will be the reflected energy.

Results for sample calculations for the reflection coefficients for various fiber structures are shown in Figs.19-20

IV. Conclusions and Recommendations

Support of this program has enabled the contractor to develop and perfect a computer program based on the scalar wave - FFT algorithm to study the propagation characteristics of guided waves in several important, practical fiber structures such as fibers with general index profiles (step index, parabolic index, ring index, etc), multi-channel fiber couplers, and fiber horns or tapers. These fiber structures may be made with commercially available fibers whose index variation may be as large as 1 - 2%.

We have implemented the adaptive coordinates and lossy outer mesh boundary schemes in our computer program. However, for most practical situations of interest in which the fiber core radius is about $25\mu\text{m}$, the cladding index is about 1 - 2% less than the core index, the spot size of the beam is less than $20\mu\text{m}$ and the free-space wavelength of the beam is larger than $0.6\mu\text{m}$, it is not necessary to implement the adaptive coordinates and lossy outer mesh boundary schemes. We also discovered that steep index gradient is not a hindrance for the program to produce accurate results as long as the scalar wave approach is still justifiable.

It is not unreasonable to ask the following question:

"Now that we have completed a beautiful program capable of producing propagation results for a variety of practical fiber structures, what can we do with it?"

The answer is "May be a lot!" Listed below are only a few of the problems that we can solve with this program:

- (1) Any single-mode or multimode weakly guiding fiber with arbitrary refractive index profile.

Our program provides the means to obtain the propagation characteristics of guided waves supported by this structure. The core of the structure may be circular, elliptical, rectangular, triangular or dumb-bell

shape with general index profile.⁸

- (2) Any fiber couplers composed of parallel strands of two or more of the above fibers. This is the only program which can provide the detailed coupling characteristics of this type of structure. Prior knowledge of coupling characteristics of a coupler is the key to successfully design and construct fiber couplers.
- (3) Any transition elements derived from the above fibers. Transition elements such as tapers, horns, or mode converters or branches can all be analyzed by our program.

Recommended Future Research

In addition to the important practical problems mentioned above that can be solved by our approach, it is worthwhile to look into the future and seek out problems of potential importance and interest. For example:

- (1) Nonlinear Fiber.

Very high intensity is achievable in fibers. One may wish to learn the propagation characteristics of waves in a fiber in which the induced nonlinearity of the medium plays a significant role. This problem may be solved by the scalar wave - FFT approach.

- (2) Mode Conversion in a Fiber Due to its statistically Varying Medium.

This problem may also be approached from the scalar wave - FFT point of view. It is known that the presence of frozen-in statistically varying medium contributes to the mode conversion phenomenon in a fiber.⁹ A systematic study of this problem will reveal the severity of this effect in changing the dispersion characteristics of beam in this fiber.

- (3) Large-size Single-Mode Fiber

The advantages of having large-size single-mode fiber are well-known. Ring-index fiber as proposed by L Eyges of RADC appears

to be a promising one. Other type of fiber whose index variation may be radially unsymmetrical, such as, a layered fiber as shown in Fig. 21 may also be promising. Research should be encouraged on this type of fibers.

(4) Polarization Preserving Fibers

One of the main features of a weakly guiding fiber which can be analyzed by the scalar wave approach is that the guiding structure is polarization insensitive. However, for several important application areas, the polarization preserving characteristic of fiber is essential. Stress induced birefringent fiber or deformed core fiber may satisfy the needed requirement. However, to achieve the targeted isolation for the two orthogonal dominant modes, the required stress is excessively high for stress induced birefringent fiber and the loss is excessive for deformed-core fiber. We propose the use of layered dielectrics as shown in Fig. 21 to produce an equivalent birefringent effect to enable proper isolation for the two orthogonal dominant modes.¹⁰ Initial indication is very promising. It would therefore be very worthwhile to pursue this research.

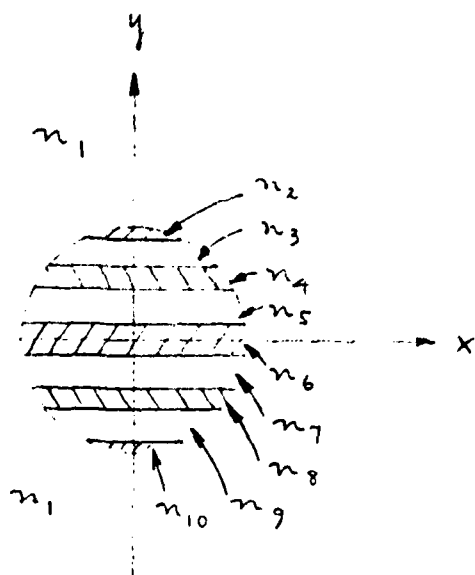


FIG. 21 Polarization Preserving Layered Fiber

Personnel:

The principal contributors of this contract have been:

C. Yeh	Senior Research Engineer
P. Barber	Research Engineer
F. Manshadi	Research Engineer

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9. M. Imai, T. Asakura and Y. Kinoshita, Opt. and Quan. Elec. 7, 95 (1975).
10. Patent Pending.

SAMPLE PROGRAM LISTINGS FOR THE CASE OF BEAM
PROPAGATION IN A STEP INDEX FIBER

CORE SIZE (DIAMETER) = $50\mu\text{m}$
CORE INDEX = 1.48
CLADDING INDEX = 1.46
INITIAL GAUSSIAN BEAM SPOT SIZE (DIAMETER) = $40\mu\text{m}$
WAVELENGTH = $0.8\mu\text{m}$

(STEP INDEX CASE, $m = 6$)

```

LEVEL 2.3 . (JUNE 78)          OS/360  QJTRAY M EXTENDED          DATE 80.19'  1.57.04

REQUESTED OPTIONS:

OPTIONS IN EFFECT:  NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTODBL(NONE)
SOURCE EBCDIC NOLIST NODECK OBJECT JJJ43 NOFORMAT GOSTAT NOXREF= ALC NOANSF NOTERM IDM FLA

C THIS PROGRAM CONTAINS:CORRECTED =18RJ4=M(JPTF1B,PRINTER,PEAK,SIZE.
C HARY4,GREYSC).
C
C PROGRAM JPTF1B(INPUT,OUTPUT,=1-E3,TAPE5=INPUT,TAPE6=OUTPUT,
C + TAPE7=FILED)
C
C INPUT PARAMETERS
C CARD 1 (15 FORMAT)
C NCASES NUMBER OF CASES TO BE READ
C CARD 2 (NAMELIST FORMAT=$DEFAULT $
C LAMBDA WAVE LENGTH (MICRONS)
C RO IYE POINT IN IRRADIANCE (MICRJNS)
C FR FIBER RADIUS (MICRONS)
C NO REFRACTIVE INDEX
C PCORP PERCENT DROP AT REF OF 1/4J
C OUTRAD OUTER RADIUS (MICRONS)
C DX WESH SPACING (MICRONS)
C NSTEPS NUMBER OF Z-STEP = 2MIN/NDZINC
C NDZINC LENGTH OF Z-STEP
C IOUT DEVICE NUMBER FOR OUTPUT
C IGREY DEVICE NUMBER FOR GREYSC JJJ2J
C PGREY IF TRUE PRINT IRRADIANCE PLTFILE AT EACH STEP
C PWALST IF TRUE WRITE 2ND ARGUMENTS AT EACH STEP
C PLTMAX IF TRUE PLOT 2ND ARGUMENTS VS DISTANCE
C PLTFLD IF TRUE PLOT PEAK INTENSITY VS DISTANCE
C PLTFILE IF TRUE PLOT FIELD IRRAD AT END OF PROPAGATION ALONG I
C WESH IF TRUE DO ABOVE PLOT AT EVERY STEP
C GRID SIZE (32, 64, OR 128)
C
C COMMON /LCM2/REFNDX(16384),SM(123),CS(129),ZTSQ(128),PQSS(1129)
C + ,AMPAY(16384),RADARY(16384)
C LEVEL 2,REFNDX,SM,CS,ZTSQ,PQSS,AMPAY,RADARY
C DIMENSION WK(256),NM(3),INVJ(2),S(32)
C DIMENSION X0(3),Y0(3),PCORP(13),=M(3),REFCFAL(3)
C COMMON /ARRAYS/V(32768)
C COMMON /PARAM/ZINC,WE SH,LAMBDA,MU,FR,NO,PCORP,OUTRAD,DX,NSTEPS.
C + NDZINC,MESHSQ,MESH502,PI,WAVEN4,2XSI,MS,MS,MF,NF,MSHPTS
C COMMON /PRINTPL/PGREY,PWALST,PLT5T,PLTMAX,LAST,IOUT,IGREY
C + ,PLTFLD,PLTFILE
C
C REAL LAMBDA,NO,MESHQ,N2
C LOGICAL PGREY,PWALST,PLTWST,PLTMAX,PLTFLD,PLTFILE,LAST,CALLPM
C
C DATA ICNTCS/1,ICNT/0/
C
C NAMELIST /DEFAULT/LAMBDA,RO,NO,PCORP,CJTRAD,DX,NSTEPS,NDZINC.
C + IOUT,IGREY,PGREY,PWALST,PLT5T,PLTMAX,PLTFLD,PLTFILE,MESH
C

```



```

LEVEL 2...J (JUNE 78)                                OS/J50  PORTMAN H EXTENDED  DATE 80.19  21.57.04

ISM 0012  READ(5,1000) NCASES
ISM 0013  WRITE(6,1000) NCASES
ISM 0014  READ(5,11)XB,YB
ISM 0015  WRITE(6,11)XB,YB
ISM 0016  READ(5,1000) NFIR
ISM 0017  WRITE(6,1000) NFIR
ISM 0018  READ(5,12)(X0(I),Y0(I),I=1,NF12)
ISM 0019  WRITE(6,12)(X0(I),Y0(I),I=1,NF12)
ISM 0020  READ(5,13)(PCORPA(K),K=1,NF13)
ISM 0021  WRITE(6,13)(PCORPA(K),K=1,NF13)
ISM 0022  READ(5,11)(FRA(K),K=1,NF18)
ISM 0023  WRITE(6,11)(FRA(K),K=1,NF18)
ISM 0024  WRITE(6,11)(FRA(K),K=1,NF18)
ISM 0025  1000 FORMAT(2X,12)
ISM 0026  11 FORMAT(1X,F4.1,1X,F4.1)
ISM 0027  12 FORMAT(1X,F4.1,1X,F4.1,1X,F4.1,1X,F12.5)
C1 13 READ(5,DEFAULT)
C

ISM 0028  MM(1)= 7
ISM 0029  MM(2)= 7
ISM 0030  MM(3)= 0
ISM 0031  I CNI=0
ISM 0032  FLAG=NJ/ABS(NO)
ISM 0033  NO=ABS(NO)
ISM 0034  IF(FLAG.LT.0.) WRITE(10UT,2050)
ISM 0035  2050 FORMAT(/,47M THE REFRACTIVE INDEX IS A CONSTANT EQUAL TO NJ)
ISM 0036  PCORP=PCORPA(1)
ISM 0037  FR=FRA(1)
ISM 0038  WRITE(10JT,DEFAULT)
ISM 0039  C

C C C
ISM 0040  CALCULATE CONSTANTS
ISM 0041  MESM2=2*MESH
ISM 0042  MESHSQ=MESH**2
ISM 0043  MSHSQ2=2*MESHSQ
ISM 0044  RN2NO=.02*PCORP/FR**2
ISM 0045  ZMIN=PI/12.*SORT(RN2NO)
ISM 0046  DZINC=ZMIN/NDZINC
ISM 0047  DXSI=DZINC/2.
ISM 0048  DZET=DX/RO
ISM 0049  WAVENM=2.*PI/LAMBDA
ISM 0050  BETAH=12.*ZMIN*DXSI/(WAVENM*NJ)*(PI/(MESH*DZET*RO))**2
ISM 0051  FICNST=1.-1./MESH)*PI
ISM 0052  XYU=MESH/2.
ISM 0053  RADNM=10J*HAD/RO)**2
ISM 0054  WNDNR=1.*MESHSQ
ISM 0055  NZ=NJ*RN2NO
ISM 0056  REFCF=N2*RU**2/2.
ISM 0057  ALPHA=2.*ZMIN/(PI*(WAVENM*NJ**4U**2))
ISM 0058  NOSQ=NJ**2
ISM 0059  KSIMUL=DXSI
ISM 0060  LAST=.FALSE.
ISM 0061  IF(FLAG.LT.0.) REFCF=0.

```

```

LEVEL 2.3.0 (JUNE 78)          MAIN          OS/350  -ATRAN M EXTENDED  DATE 80.199/  .57.04

ISN 0063      CALL PR=PGREY,OR,PWAIST,OR,P,INSTR,J4,2,TRAX,OR,PLTFLD,OR,PLT=LE
C
C      WRITE THE IMPORTANT CALCULATED PARAMETERS
C
ISN 0064      WRITE(1000) ZMIN,DZINC,MN2VJ,ALPHA
ISN 0065      FORMAT(1.8M ZMIN = ,F10.4,1X,714(-40VS,/.9H DZINC = ,F10.4,1X,
+ 7HMICRONS,/.9H MN2NO = ,F10.3,1X,1HMICRONS*(-2),/,
+ 9H ALPHA = ,F10.5,/)
C
C      CALCULATE NECESSARY ARRAYS
C
MP=6
DO 800 K=1,NFIB
  X0(K)=X0(K)*XYD
  Y0(K)=Y0(K)*XYD
  MX=2*MP
  REFCFA(K)=Y0*.02*PCORPA(K)*((30/2*Z(K))*MX/2.
CONTINUE
DO 100 K=1,MESH
  RK=K-1
  ARGFCNST*RK
  CS(K)=COS(ARG)
  SN(K)=SIN(ARG)
  ZTSQ(K)=((RK-XYD)*DZET)**2
  POSQ(K)=((RK-XYD+.5)**2
CONTINUE
100
C
C      SET UP REFRACTIVE INDEX ARRAY
C
M=0
DO 120 J=1,MESH
  DC 120 I=1,MESH
  N=M+1
  TMPNDX=0.
  DO 110 K=1,NFIB
    Z1=((J-1-Y0(K))*DZET)**2
    Z2=((I-1-X0(K))*DZET)**2
    RAD=Z1+Z2
    TMPNDX=AMAX1(TMPNDX,(NO-REFCFA(K))*((NAD**MPI)))
    REF=NUS(1-D.01*PCORPA(K))
    REFNDX(N)=AMAX1(TMPNDX,REF)
CONTINUE
110
C
C      WRITE(1100,*) REFNDX(M)
C
DO 120 I=1,MESH
  IF MESH.NE.128) GO TO 10
  MS=MESH/4+1
  MF=MS+MESH/2-1
  NS=MS
  NF=MF
  GO TO 40
CONTINUE
10
MS=1
MF=MESH
ISN 0081
ISN 0082
ISN 0083
ISN 0084
ISN 0085
ISN 0086
ISN 0087
ISN 0088
ISN 0089
ISN 0090
ISN 0091
ISN 0092
ISN 0093
ISN 0094
ISN 0095
ISN 0096
ISN 0097
ISN 0098
ISN 0099
ISN 0100
ISN 0101
ISN 0102
ISN 0103
ISN 0104
ISN 0105

```


LEVEL 2.3.	(JUNE 78)	MAIN	05/360	JATRAM EXTENDED	DATE 80-199	1.57.04
ISN 0150		V(K)=VR*AR-VI*AI				OPTF18
ISN 0151		V(KPI)=VI*AR+VR*AI				OPTF18
ISN 0152		CONTINUE				OPTF18
	160					OPTF18
	C	DO TRANS=JRM				OPTF18
	C	CALL HARM(V*MM,INV,S.1,IFERR)				OPTF18
ISN 0153		SOLVE FIRST ORDER ODE				OPTF18
	C					OPTF18
	C					OPTF18
ISN 0154		K=-1				OPTF18
ISN 0155		DO 190 J=1,MESH				OPTF18
ISN 0156		PHI1=SETAMT*POSQ(J)				OPTF18
ISN 0157		DO 180 I=1,MESH				OPTF18
ISN 0158		K=K+2				OPTF18
ISN 0159		KPI=K+1				OPTF18
ISN 0160		PHI2=SETAMT*POSQ(I)				OPTF18
ISN 0161		VR=V(I)				OPTF18
ISN 0162		VI=V(KPI)				OPTF18
ISN 0163		ANG2=-(PHI1*PHI2)				OPTF18
ISN 0164		CANG=CDS(ANG)				OPTF18
ISN 0165		SANG=SLY(ANG)				OPTF18
ISN 0166		V(K)=IVR*CANG-VIS*ANG				OPTF18
ISN 0167		V(KPI)=-(VR*SANG+VI*CANG)				OPTF18
ISN 0168		CONTINUE				OPTF18
	180					OPTF18
	C	DO INVERSE TRANSFORM				OPTF18
	C					OPTF18
ISN 0169		CALL HARM(V*MM,INV,S.-1,IFERR)				OPTF18
	C	RECONDITION V BECAUSE OF TRANSFORM				OPTF18
	C					OPTF18
	C					OPTF18
ISN 0170		K=-1				OPTF18
ISN 0171		DO 200 J=1,MESH				OPTF18
ISN 0172		SAJ=SN(J)				OPTF18
ISN 0173		CSJ=CS(J)				OPTF18
ISN 0174		RZ1=(1-J)*YO-1)*OZET)*2				OPTF18
ISN 0175		CLOSSE=2				OPTF18
ISN 0176		DO 200 I=1,MESH				OPTF18
ISN 0177		BZ2=(1-I)*YO-1)*OZET)*2				OPTF18
ISN 0178		BRAD=BZ1*BZ2				OPTF18
ISN 0179		CM=1				OPTF18
ISN 0180		IF(BZ1*DO.GT.HADNRM) CM=EXP(-C.USS*(HADNRM-RADNRM))				OPTF18
ISN 0181		K=K+2				OPTF18
ISN 0182		KPI=K+1				OPTF18
ISN 0183		SN=SN(I)				OPTF18
ISN 0184		CS=CS(I)				OPTF18
ISN 0185		AR=CSJ*CS1-SNJ*SN1				OPTF18
ISN 0186		AI=CSJ*SN1+SNJ*CS1				OPTF18
ISN 0187		VR=V(K)*CM				OPTF18
ISN 0188		VI=V(KPI)*CM				OPTF18
ISN 0189		V(K)=VR*AR-VI*AI				OPTF18
ISN 0190		V(KPI)=VR*AI+VI*AR				OPTF18
ISN 0191		CONTINUE				OPTF18
ISN 0192						OPTF18

```

LEVEL 2.3.0 (JUNE 78)      MAIN      QS/350  =QJTRAN H EXTENDED      DATE 80.199/21.57.04

C      NOW INCLUDE EITHER FULL STEP U3 1A-F STEP REFRACTIVE
C      INDEX EFFECTS DEPENDING ON #HERE IN THE PATH YOU ARE
C
ISM 0193      IF(ICNT.EQ.NSTEPS) XSIMUL=DXSIN
ISM 0194      K=1
ISM 0195      CH=WAVERN*ZMIN*XSIMUL/(2.*NCJ)
ISM 0196      DO 220 M=1,MESHSD
ISM 0197      ARG=RE-NDX(M)*CH
ISM 0198      AR=CDS(ARG)
ISM 0199      AT=SIN(ARG)
ISM 0200      KSK+2
ISM 0201      KPI=K+1
ISM 0202      V=V(K)
ISM 0203      V=V(KPI)
ISM 0204      V(K)=VR*AR-VI*AT
ISM 0205      V(KPI)=V*AT+VI*AR
ISM 0206      CONTINUE
ISM 0207      IF(ICNT.EQ.NSTEPS) LAST=.TRUE.
ISM 0208      IF(CALL3) CALL PRINTER(ICNT)
ISM 0210      220 CONTINUE
ISM 0212
C
C      CALCULATE IRRADIANCE PATTERN AND PRINT
C
ISM 0213      IF(PGREY) GO TO 30
ISM 0214      N=0
ISM 0215      DO 240 K=1,MESHSD2.2
ISM 0216      KPI=K+1
ISM 0217      VR=V(K)
ISM 0218      VI=V(KPI)
ISM 0219      M=M+1
ISM 0220      RADARY(4)=VR**2+VI**2
ISM 0221      240 CALL GREYSCALE(1,REY,10,RADARY,MESH4,MESH4,MS,MF,1,NS,MF,1,0..0..,
ISM 0222      * 10MIRADIANCE,10)
ISM 0223
ISM 0224      30 CONTINUE
ISM 0225      ICNTCS=ICNTCS+1
ISM 0226      IF(ICNTCS.LE.NCASES) GO TO 1
ISM 0227      STOP
END

*OPTIONS IN EFFECT*NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTOOBL(NONE)
*OPTIONS IN EFFECT*SOURCE EBCDIC NOLIST NODECK OBJECT UNMAP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IRM F1A
*STATISTICS* SOURCE STATEMENTS = 226, PROGRAM SIZE = 5316, SUBPROGRAM NAME = MAIN
*STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
68K BYTES OF CORE NOT USED

```

LEVEL 2.3.3 (JUNE 79) OS/360 J3711AN H EXTENDED DATE 80.195 1.57.07

REQUESTED OPTIONS:

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTODBL(NONE) ALC NOANSF NOTERM IBM FLA
SOURCE EBCDIC NOLIST NOCHECK OBJECT NJ4AP NOFORMAT GOSTMT NOXREF

```

      SUBROUTINE PEAK(VMAX)
      COMMON /LCM2/REFNDX(16384),JN(123),CS(129),ZTSO(128),POS(128)
      * .ANPARY(16384),RADARY(16384)
      LEVEL 2,REFNDX,SN,CS,ZTSO,POSU,ANPARY,RADARY
      COMMON /ARRAYS/V(32768)
      COMMON /ZARAY/DZINC,MESH,LAYJDA,Q,F,N),PCOMP,OUTRAD,DX,NSTEPS,
      * NOZINC,MESHQ,MHSO2,PI,WAVENM,XS,I,MS,NS,MF,NF,MSHPTS
      L=0
      SUML=0.
      SUMR=0.
      VMAX=0.
      DO 10 K=1,MSHS32*2
      VR=V(K)
      L=L+1
      KPI=K+1
      VI=V(KPI)
      VRAD=VI**2+VR**2
      IF(K.LE.MESHQ) SUML=SUML+VRAD
      IF(K.GE.MESHQ) SUMR=SUMR+VRAD
      VMAX=AMAX1(VMAX,VRAD)
      RADARY(L)=VRAD
      CONTINUE
      TOT=SUML+SUMR
      SUML=SUML/TOT
      SUMR=SUMR/TOT
      WRITE(6,2000) SUML,SUMR
      2000 FORMAT(1X,7H$UML = ,E14.7,3X,7H$JMR = ,E14.7,/)
      RETURN
      END
    10
  
```

112K BYTES OF CORE NOT USED

***** END OF COMPILATION *****

***** NO DIAGNOSTICS GENERATED *****

***** SOURCE STATEMENTS = 28, PROGRAM SIZE = 540, SUBPROGRAM NAME = PEAK *****

***** OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTODBL(NONE) *****

***** OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NOCHECK OBJECT NJ4AP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IBM FL *****

***** STATISTICS *****

***** STATISTICS *****

```

LEVEL 2.3. (JUNE 78)
REQUESTED OPTIONS:
OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE)
SOURCE EBCDIC NOLIST NODECK OBJECT VJ4AP NOFORMAT GOSTMT NOKHEF ALC NOANSF NOTERM IR4 FL/

SURROUTINE SIZE(ARRAY,MESH,MESH,XO,YO,X2,Y2)
DIMENSION ARRAY(MESH,MESH)
MID=MESH/2+1
SUM1=0.
SUM2=0.
SUM3=0.
SUM4=0.
SUM5=0.
DO 10 N=1,MESH
RN=N-MID
RNSQ=RN**2
DO 10 M=1,MESH
RMSQ=RMSQ+RNSQ
RMSQ=RMSQ**2
ARRAYN=ARRAY(M,N)
SUM1=SUM1+ARRAYN
SUM2=SUM2+RMSQ
SUM3=SUM3+RMSQ*ARRAYN
SUM4=SUM4+SUM3*ARRAYN
SUM5=SUM5+ARRAYN
CONTINUE
SNORM=1./SUM5
XO=SNORM*SUM2
YO=SNORM*SUM1
X2=SNORM*SUM4-XO**2
Y2=SNORM*SUM3-YO**2
X2=SQRT(X2)
Y2=SQRT(Y2)
RETURN
END

10

OPTIONS IN EFFECT:NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODBL(NONE)
OPTIONS IN EFFECT:SOURCE EBCDIC NOLIST NODECK OBJECT VJ4AP NOFORMAT GOSTMT NOKHEF ALC NOANSF NOTERM IR4 FLA
*STATISTICS* SOURCE STATEMENTS = 30. PROGRAM SIZE = 786. SUBPROGRAM NAME = SIZE
*STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
112K BYTES OF CORE NOT USED

```


15N 0003
15N 0004
15N 0005
15N 0006
15N 0007
15N 0008
15N 0009
15N 0010
15N 0011
15N 0012
15N 0013
15N 0014
15N 0015
15N 0016
15N 0017
15N 0019

LEVEL 2.	J (JUNE 78)	HARM	05/360 J3TRAN H EXTENDED	DATE 80.19 21.57.09
ISN 0020		FN = NX		HARM1070
ISN 0021		DD 19 I = 1, NX		HARM1080
ISN 0022		A(2+I-1) = A(2+I-1)/FN		HARM1090
ISN 0023		19 A(2+I) = -A(2+I)/FN		HARM1100
ISN 0024		20 NP(1) = N1+2		HARM1110
ISN 0025		NP(2) = NP(1)+N2		HARM1120
ISN 0026		NP(3) = NP(2)+N3		HARM1130
ISN 0027		DD 250 I = 1, N3		HARM1140
ISN 0028		IL = NP(3) - NP(10)		HARM1150
ISN 0029		IL1 = IL+1		HARM1160
ISN 0030		MI = M(10)		HARM1170
ISN 0031		IF (N1)250,250,30		HARM1180
ISN 0032		30 IDIF=NP(10)		HARM1190
ISN 0033		KBIT=NP(10)		HARM1200
ISN 0034		MEV = 2*(MI/2)		HARM1210
ISN 0035		IF (N1 - MEV)60,60,40		HARM1220
	C			HARM1230
	C	M IS ODD, DD L=1 CASE		HARM1240
ISN 0036		40 KBIT=CBIT/2		HARM1250
ISN 0037		KL=KBIT-2		HARM1260
ISN 0038		DD 50 I=1, IL1, IDIF		HARM1270
ISN 0039		KLAST=KL+1		HARM1280
ISN 0040		DD 50 K=(, KLAST, 2		HARM1290
ISN 0041		KD=K+KBIT		HARM1300
	C			HARM1310
	C	DD ONE STEP WITH L=1, J=0		HARM1320
	C	A(K)=A(K)+A(KD)		HARM1330
	C	A(KD)=A(K)-A(KD)		HARM1340
	C			HARM1350
ISN 0042		T=A(KD)		HARM1360
ISN 0043		A(KD)=A(K)-T		HARM1370
ISN 0044		A(K)=A(K)+T		HARM1380
ISN 0045		T=A(KD+1)		HARM1390
ISN 0046		A(KD+1)=A(K+1)-T		HARM1400
ISN 0047		50 A(K+1)=A(K+1)+T		HARM1410
ISN 0048		IF (N1 - 1)250,250,52		HARM1420
ISN 0049		52 LFIRST = 3		HARM1430
	C			HARM1440
	C	DEF - JLAST = 2*(L-2) - 1		HARM1450
ISN 0050		JLAST=1		HARM1460
ISN 0051		GO TO 70		HARM1470
	C			HARM1480
	C	M IS EVEN		HARM1490
ISN 0052		60 LFIRST = 2		HARM1500
ISN 0053		JLAST=0		HARM1510
ISN 0054		70 DD 240 L=LFIRST, MI, 2		HARM1520
ISN 0055		JJOF=KBIT		HARM1530
ISN 0056		KBIT=KBIT/4		HARM1540
ISN 0057		KL=KBIT-2		HARM1550
	C			HARM1560
	C	DD FJH J=0		HARM1570
ISN 0058		DD 80 I=1, IL1, IDIF		HARM1580
ISN 0059		KLAST=I+KL		HARM1590
ISN 0060		DD 80 K=I, KLAST, 2		HARM1600

LEVEL 2.	(JUNE 78)	HARM	OS/350	JRTIAN H EXTENDED	DATE	80-15	21.57.09
ISM 0061		K1=K+KBIT				HARM1610	
ISM 0062		K2=K1+KBIT				HARM1620	
ISM 0063		K3=K2+KBIT				HARM1630	
		DO TWO STEPS WITH J=0				HARM1640	
		A(K)=A(K)+A(K2)				HARM1650	
		A(K2)=A(K)-A(K2)				HARM1660	
		A(K1)=A(K1)+A(K3)				HARM1670	
		A(K3)=A(K1)-A(K3)				HARM1680	
		A(K)=A(K)+A(K1)				HARM1690	
		A(K1)=A(K)-A(K1)				HARM1700	
		A(K2)=A(K2)+A(K3)+1				HARM1710	
		A(K3)=A(K2)-A(K3)+1				HARM1720	
		T=A(K2)				HARM1730	
		A(K2)=A(K)-T				HARM1740	
		A(K)=A(K)+T				HARM1750	
		T=A(K2+1)				HARM1760	
		A(K2+1)=A(K+1)-T				HARM1770	
		A(K+1)=A(K+1)+T				HARM1780	
		T=A(K3)				HARM1790	
		A(K3)=A(K1)-T				HARM1800	
		A(K1)=A(K1)+T				HARM1810	
		T=A(K3+1)				HARM1820	
		A(K3+1)=A(K+1)+T				HARM1830	
		T=A(K1)				HARM1840	
		A(K1)=A(K)-T				HARM1850	
		A(K)=A(K)+T				HARM1860	
		T=A(K1+1)				HARM1870	
		A(K1+1)=A(K+1)-T				HARM1880	
		A(K+1)=A(K+1)+T				HARM1890	
		R=A(K3+1)				HARM1900	
		T=A(K3)				HARM1910	
		A(K3)=A(K2)-R				HARM1920	
		A(K2)=A(K2)+R				HARM1930	
		A(K3+1)=A(K2+1)-T				HARM1940	
		A(K2+1)=A(K2+1)+T				HARM1950	
		IF (JLAST) 235.235.02				HARM1960	
		80 JJ=JJ+1				HARM1970	
		82 JJ=JJ+1				HARM1980	
		DO FOR J=1				HARM1990	
		JLAST=1L+JJ				HARM2000	
		DO 85 I=JJ, JLAST, IDIF				HARM2010	
		KLAST=K+1				HARM2020	
		DO 83 K=1, KLAST, 2				HARM2030	
		K1=K+KBIT				HARM2040	
		K2=K+KBIT				HARM2050	
		K3=K+KBIT				HARM2060	
						HARM2070	
						HARM2080	
						HARM2090	
						HARM2100	
						HARM2110	
						HARM2120	
						HARM2130	
						HARM2140	

LEVEL 2.3. - (JUNE 78)	HARM	DS/350	J413AN M EXTENDED	DATE 80.199	1.57.09
15N 0097	C	LETTING W=(1+I)/ROOT2.W3=(-1+I)/ROOT2.W2=I.		HARM2150	
15N 0098	C	A(K1)=A(K1)+A(K2)*I		HARM2160	
15N 0099	C	A(K2)=A(K1)-A(K2)*I		HARM2170	
15N 0100	C	A(K1)=A(K1)+W*A(K3)*W3		HARM2180	
15N 0101	C	A(K3)=A(K1)+W-A(K3)*W3		HARM2190	
15N 0102	C	A(K1)=A(K1)+A(K1)		HARM2200	
	C	A(K2)=A(K1)-A(K1)		HARM2210	
	C	A(K3)=A(K2)+A(K3)*I		HARM2220	
	C	A(K1)=A(K2)-A(K3)*I		HARM2230	
	C	R = -A(K2)		HARM2240	
	C	T = A(K2)		HARM2250	
	C	A(K1) = A(K1)-R		HARM2260	
	C	A(K2) = A(K1)+R		HARM2270	
	C	A(K3+1)=A(K1)-T		HARM2280	
	C	A(K1)=A(K1)+T		HARM2290	
	C	AW=A(K1)-A(K1+1)		HARM2300	
	C	AW1 = A(K1+1)+A(K1)		HARM2310	
	C	R = A(K3)-A(K3+1)		HARM2320	
	C	T = A(K3)-A(K3+1)		HARM2330	
	C	A(K3)=(AW-R)/ROOT2		HARM2340	
	C	A(K3+1)=(AW1-T)/ROOT2		HARM2350	
	C	A(K1)=(AW+R)/ROOT2		HARM2360	
	C	A(K1+1)=(AW1+T)/ROOT2		HARM2370	
	C	T = A(K1)		HARM2380	
	C	A(K1)=A(K1)-T		HARM2390	
	C	A(K2)=A(K1)+T		HARM2400	
	C	T=A(K1+1)		HARM2410	
	C	A(K1+1)=A(K1+1)-T		HARM2420	
	C	A(K1)=A(K1)+T		HARM2430	
	C	R = -A(K3+1)		HARM2440	
	C	T = A(K3)		HARM2450	
	C	A(K3)=A(K2)-R		HARM2460	
	C	A(K2)=A(K2)+R		HARM2470	
	C	A(K3+1)=A(K2+1)-T		HARM2480	
	C	A(K2+1)=A(K2+1)+T		HARM2490	
	C	IF(LAST-I) 235,235,90		HARM2500	
85		JJ= JJ + JJ0F		HARM2510	
90		JJ= JJ + JJ0F		HARM2520	
	C	NOW DO THE REMAINING J'S		HARM2530	
	C	DO 230 J=2, LAST		HARM2540	
	C	FETCH W'S		HARM2550	
	C	DEF W=0*INV(J), W2=W**2, WJ=W**3		HARM2560	
96		I=INV(J+1)		HARM2570	
98		IC=NI-I		HARM2580	
	C	W(1)=S(IC)		HARM2590	
	C	W(2)=S(I)		HARM2600	
	C	I2=2*I		HARM2610	
	C	IC=NI-I2		HARM2620	
	C	IF(I2C)I20, I10, I00		HARM2630	
	C			HARM2640	
	C			HARM2650	
	C			HARM2660	
	C			HARM2670	
	C			HARM2680	
	C			HARM2690	
	C			HARM2700	
	C			HARM2710	
	C			HARM2720	
	C			HARM2730	
	C			HARM2740	
	C			HARM2750	
	C			HARM2760	
	C			HARM2770	
	C			HARM2780	
	C			HARM2790	
	C			HARM2800	
	C			HARM2810	
	C			HARM2820	
	C			HARM2830	
	C			HARM2840	
	C			HARM2850	
	C			HARM2860	
	C			HARM2870	
	C			HARM2880	
	C			HARM2890	
	C			HARM2900	
	C			HARM2910	
	C			HARM2920	
	C			HARM2930	
	C			HARM2940	
	C			HARM2950	
	C			HARM2960	
	C			HARM2970	
	C			HARM2980	
	C			HARM2990	
	C			HARM3000	

LEVEL 20	(JUNE 78)	HARM	OS/350	OUTRAN H EXTENDED	DATE 80.19	21.57.09
ISN 0133	C	201 IS IN FIRST QUADRANT				HARM2690
ISN 0134		100 W2(1)=S(12C)				HARM2700
ISN 0135		W2(2)=S(12)				HARM2710
ISN 0136		GO TO 130				HARM2720
ISN 0137		110 W2(1)=0.				HARM2730
ISN 0138		W2(2)=1.				HARM2740
		GO TO 130				HARM2750
ISN 0139	C	201 IS IN SECOND QUADRANT				HARM2760
ISN 0140		120 12CC = 12C+NT				HARM2770
ISN 0141		12C = 12C				HARM2780
ISN 0142		W2(1)=S(12C)				HARM2790
ISN 0143		W2(2)=S(12CC)				HARM2800
ISN 0144		130 13C=12				HARM2810
ISN 0145		13C=NT+13				HARM2820
		IF(13C)160.150.140				HARM2830
ISN 0146	C	13 IN FIRST QUADRANT				HARM2840
ISN 0147		140 W3(1)=S(13C)				HARM2850
ISN 0148		W3(2)=S(13)				HARM2860
ISN 0149		GO TO 200				HARM2870
ISN 0150		150 W3(1)=0.				HARM2880
ISN 0151		W3(2)=1.				HARM2890
		GO TO 200				HARM2900
ISN 0152	C	160 13CC=13C+NT				HARM2910
ISN 0153		IF(13CC)190.180.170				HARM2920
ISN 0154	C	13 IN SECOND QUADRANT				HARM2930
ISN 0155		170 13C=13C				HARM2940
ISN 0156		W3(1)=S(13C)				HARM2950
ISN 0157		W3(2)=S(13CC)				HARM2960
ISN 0158		GO TO 200				HARM2970
ISN 0159		180 W3(1)=1.				HARM2980
ISN 0160		W3(2)=0.				HARM2990
		GO TO 200				HARM3000
ISN 0161	C	301 IN THIRD QUADRANT				HARM3010
ISN 0162		190 13CC=NT+13C				HARM3020
ISN 0163		13CC = 13CC				HARM3030
ISN 0164		W3(1)=S(13CC)				HARM3040
ISN 0165		W3(2)=S(13CC)				HARM3050
ISN 0166		1LAST=1L+JJ				HARM3060
ISN 0167		DO 220 1=JJ.1LAST.1DIF				HARM3070
ISN 0168		1LAST=KL+I				HARM3080
ISN 0169		DO 220 K=1.1LAST.2				HARM3090
ISN 0170		K1=K+K3IT				HARM3100
ISN 0171		K2=K1+K3IT				HARM3110
		K3=K2+K3IT				HARM3120
ISN 0172	C	DO TWO STEPS WITH J NOT 0				HARM3130
ISN 0173		A(KJ)=A(KJ)+A(K2)*W2				HARM3140
ISN 0174		A(K2)=A(K2)-A(K2)*W2				HARM3150
ISN 0175		A(K1)=A(K1)+W2*A(K3)*W3				HARM3160
ISN 0176						HARM3170
ISN 0177						HARM3180
ISN 0178						HARM3190
ISN 0179						HARM3200
ISN 0180						HARM3210
ISN 0181						HARM3220

LEVEL 2.3.0 (JUNE 78)	HARM	DS/360	33TRAN M EXTENDED	DATE 80.195	1.57.09
ISN 0172	C	A(K3)=A(K1)*W-A(K3)*W3		HARM3230	
ISN 0173	C	A(K1)=A(K1)+A(K1)		HARM3240	
ISN 0174	C	A(K1)=A(K1)-A(K1)		HARM3250	
ISN 0175	C	A(K2)=A(K2)+A(K3)*I		HARM3260	
ISN 0176	C	A(K3)=A(K2)-A(K3)*I		HARM3270	
ISN 0177	C	R=A(K2)*W2(1)-A(K2+1)*W2(2)		HARM3280	
		T=A(K2)*W2(2)+A(K2+1)*W2(1)		HARM3290	
		A(K2)=A(K)-R		HARM3300	
		A(K)=A(K)+R		HARM3310	
		A(K2+1)=A(K+1)-T		HARM3320	
		A(K+1)=A(K+1)+T		HARM3330	
		R=A(K3)*W3(1)-A(K3+1)*W3(2)		HARM3340	
		T=A(K3)*W3(2)+A(K3+1)*W3(1)		HARM3350	
		AWR=A(K1)*W(1)-A(K1+1)*W(2)		HARM3360	
		AWI=A(K1)*W(2)+A(K1+1)*W(1)		HARM3370	
		A(K3)=A(K2)-R		HARM3380	
		A(K3+1)=A(K1)-T		HARM3390	
		A(K1)=AWR+R		HARM3400	
		T=A(K1)-AWI+T		HARM3410	
		A(K1)=A(K1)-T		HARM3420	
		A(K)=A(K)+T		HARM3430	
		T=A(K+1)		HARM3440	
		A(K+1)=A(K+1)+T		HARM3450	
		R=A(K3+1)		HARM3460	
		A(K1+1)=A(K+1)-T		HARM3470	
		A(K+1)=A(K+1)+T		HARM3480	
		T=A(K3)		HARM3490	
		A(K3)=A(K2)-R		HARM3500	
		A(K2)=A(K2)+R		HARM3510	
		A(K3+1)=A(K2+1)-T		HARM3520	
		A(K2+1)=A(K2+1)+T		HARM3530	
220	C	END 3F I AND K LOOPS		HARM3540	
				HARM3550	
				HARM3560	
				HARM3570	
				HARM3580	
				HARM3590	
				HARM3600	
				HARM3610	
				HARM3620	
				HARM3630	
				HARM3640	
				HARM3650	
				HARM3660	
				HARM3670	
				HARM3680	
				HARM3690	
				HARM3700	
				HARM3710	
				HARM3720	
				HARM3730	
				HARM3740	
				HARM3750	
				HARM3760	
ISN 0198	C	JJ=JJDIF+JJ			
		END 3F J-LOOP			
ISN 0199	C	JLAST=JLAST+3			
ISN 0200	C	CONTINUE			
		END 3F L LOOP			
ISN 0201	C	CONTINUE			
		END 3F ID LOOP			
ISN 0202	C	WE NOW HAVE THE COMPLEX FOURIER SUMS BUT THEIR ADDRESSES ARE			
ISN 0203	C	BIT-REVERSED. THE FOLLOWING ROUTINE PUTS THEM IN ORDER			
ISN 0204	C	NTSQ=NT			
		NT=NT-4			
		IF(MJMT) 370,360,360			
ISN 0205	C	M3 G4. ON EQ. MT			
		M3 G03=1			

LEVEL 2.3.0 (JUNE 78)	HARM	05/300	JATTAN H EXTENDED	DATE 80.194	.57.00
ISN 0206	N3VNT=N3/NT			HARM3770	
ISN 0207	MINN3=NT			HARM3780	
ISN 0208	GO TO 380			HARM3790	
				HARM3800	
ISN 0209	M3 LESS THAN MT			HARM3810	
ISN 0210	IG03=2			HARM3820	
ISN 0211	N3VNT=1			HARM3830	
ISN 0212	NTVN3=NT/N3			HARM3840	
ISN 0213	MINN3=N3			HARM3850	
ISN 0214	JJ03 = NT50/N3			HARM3860	
ISN 0215	M2MT=M2-MT			HARM3870	
	450 IF (M2MT)470.460.460			HARM3880	
				HARM3890	
ISN 0216	M2 GR. OR EQ. MT			HARM3900	
ISN 0217	IG02=1			HARM3910	
ISN 0218	N2VNT=N2/NT			HARM3920	
ISN 0219	MINN2=NT			HARM3930	
	GO TO 480			HARM3940	
				HARM3950	
ISN 0220	M2 LESS THAN MT			HARM3960	
ISN 0221	IG02 = 2			HARM3970	
ISN 0222	N2VNT=1			HARM3980	
ISN 0223	NTVN2=NT/N2			HARM3990	
ISN 0224	MINN2=N2			HARM4000	
ISN 0225	JJ02=NT50/N2			HARM4010	
ISN 0226	M1MT=M1-MT			HARM4020	
	550 IF (M1MT)570.560.560			HARM4030	
				HARM4040	
ISN 0227	M1 GR. OR EQ. MT			HARM4050	
ISN 0228	IG01=1			HARM4060	
ISN 0229	N1VNT=N1/NT			HARM4070	
ISN 0230	MINN1=NT			HARM4080	
	GO TO 580			HARM4090	
				HARM4100	
ISN 0231	M1 LESS THAN MT			HARM4110	
ISN 0232	IG01=2			HARM4120	
ISN 0233	N1VNT=1			HARM4130	
ISN 0234	NTVN1=NT/N1			HARM4140	
ISN 0235	MINN1=NT			HARM4150	
ISN 0236	JJ01=NT50/N1			HARM4160	
ISN 0237	JJ3=1			HARM4170	
ISN 0238	DD 890 JJ3=1.N3VNT			HARM4180	
ISN 0239	IP3=INV(JJ3)			HARM4190	
ISN 0240	DD 870 JJ3=1.M1NN3			HARM4200	
ISN 0241	GO TO (610.620).IG03			HARM4210	
ISN 0242	IP3=INV(JJ3).N3VNT			HARM4220	
ISN 0243	GO TO 630			HARM4230	
ISN 0244	IP3=INV(JJ3).NTVN3			HARM4240	
ISN 0245	IP3=INV(JJ3).N2			HARM4250	
ISN 0246	JJ2=1			HARM4260	
ISN 0247	DD 870 JJ2=1.N2VNT			HARM4270	
ISN 0248	IP2=INV(JJ2).I3			HARM4280	
ISN 0249	DD 850 JJ2=1.M1NN2			HARM4290	
				HARM4300	

LEVEL 2-3-	(JUNE 78)	HARM	OS/350	JSTRAN H EXTENDED	DATE 80-199, -57.09
15N 0230	GO TO (710,720),IG02				HARM4310
15N 0231	IP2=INV(JP2)*N2VNT				HARM4320
15N 0232	GO TO 730				HARM4330
15N 0233	720 IP2=INV(JP2)/N2VNT				HARM4340
15N 0234	730 I2=((JP2+IP2)*N1				HARM4350
15N 0235	800 JJI=1				HARM4360
15N 0236	800 850 J2=1,N1VNT				HARM4370
15N 0237	IPPI=INV(JJI)+12				HARM4380
15N 0238	80 850 J1=1,MINNI				HARM4390
15N 0239	GO TO (810,820),IG01				HARM4400
15N 0240	810 IP1=INV(JP1)*N1VNT				HARM4410
15N 0241	GO TO 830				HARM4420
15N 0242	820 IP1=INV(JP1)/N1VNT				HARM4430
15N 0243	830 I2=((IP1+IP1)+1				HARM4440
15N 0244	840 IF (J-1) 840,850,850				HARM4450
15N 0245	840 T=A(I)				HARM4460
15N 0246	A(I)=A(J)				HARM4470
15N 0247	A(J)=T				HARM4480
15N 0248	A(I+1)=A(J+1)				HARM4490
15N 0249	A(J+1)=T				HARM4500
15N 0250	850 J=J2				HARM4510
15N 0251	860 JJI=JJI+JJ1				HARM4520
15N 0252	END OF JP1 AND JP2				HARM4530
15N 0273	870 JJ2=JJ2+JJ2				HARM4540
15N 0274	END OF JP2 AND JP3 LOOPS				HARM4550
15N 0275	880 JJ3 = JJ3+JJ03				HARM4560
15N 0276	END OF JPP3 LOOP				HARM4570
15N 0277	890 IF (IFSET) 891,895,895				HARM4580
15N 0278	891 DO 892 1 = 1,NX				HARM4590
	892 A(2+1) = -A(2+1)				HARM4600
	895 RETURN				HARM4610
					HARM4620
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					HARM6990
					HARM7000

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LEVEL 2.3-- (JUNE 78)      HARM      OS/350      JSTRAN H EXTENDED      DATE 80.199.  .57.04
ISN 0288      S(JDIF)=SIN(THETA)
ISN 0289      DO 950 L=2,MT
ISN 0290      THETA=THETA/2.
ISN 0291      JSTEP2=JSTEP
ISN 0292      JSTEP=JDIF
ISN 0293      JDIF=JSTEP/2
ISN 0294      S(JDIF)=SIN(THETA)
ISN 0295      JCI=NT-JDIF
ISN 0296      S(JCI)=COS(THETA)
ISN 0297      JLAST=VT-JSTEP2
ISN 0298      IF(JLAST - JSTEP) 950,920,920
ISN 0299      DO 940 J=JSTEP,JLAST,JSTEP
ISN 0300      JC=NT-J
ISN 0301      JD=J-JDIF
ISN 0302      940 S(JD)=S(J)*S(JCI)+S(JDIF)*S(JC)
ISN 0303      950 CONTINUE
C
C      SET UP INV(J) TABLE
C
ISN 0304      960 MLEXP=NTV2
C
C      MLEXP=2*(MT-L). FOR L=1
C      MLEXP=1
C
ISN 0305      LMLEXP=2*(L-1). FOR L=1
C
ISN 0306      INV(1)=0
ISN 0307      DO 970 L=1,MT
ISN 0308      INV(LMLEXP+1) = MLEXP
ISN 0309      DO 970 J=2,LVLEXP
ISN 0310      JJ=J+LMLEXP
ISN 0311      INV(JJ)=INV(J)+MLEXP
ISN 0312      MLEXP=MLEXP/2
ISN 0313      970 MLEXP=LMLEXP/2
ISN 0314      980 IF(IFSET)12,895,12
ISN 0315      982 IF(IFSET)12,895,12
C      END

*OPTIONS IN EFFECT*NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTOOBL(NONE)
*OPTIONS IN EFFECT*SOURCE EBCDIC NOLIST NODECK OBJECT UNMAP NJFORMAT GOSTMT NORREF ALC NUANSF NOTERN IOM FL/
*STATISTICS* SOURCE STATEMENTS = 314, PROGRAM SIZE = 5206, SUBPROGRAM NAME = HARM
*STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
56K BYTES OF CORE NOT USED

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[illegible]

LEVEL 2.3.0 (JUNE 78)	GREYSC	OS/360	ATTN H EXTENDED	DATE 80.199, .57.13
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15N 0044 DO 400 I=IMIN,IMAX,1DEL
15N 0045 DO 250 J=1,3
15N 0046 DO 250 J=2,NLASTL
15N 0047 250 LINE(J,1)=1BLANK
15N 0048 LEV=1
15N 0049 J=2
15N 0050 DO 300 JJ=JMIN,JMAX,1DEL
15N 0051 J=JJ+1
15N 0052 L=(AMAT(I,JJ)-ANN)/(AMX-ANN)*FLJAT(NLEVEL)+1.
15N 0053 L=MAX(1,L)
15N 0054 L=MIN(1,NLEVEL,L)
15N 0055 LEV=MAX(1,LEV,LEVEL(L))
15N 0056 LEVNOW=LEVEL(L)
15N 0057 DO 300 K=1,LEVNOW
15N 0058 LINE(J,K)=1CHARS(L,K)
15N 0059 300 CONTINUE

C
C FIND LAST PRINT POSITION
C
15N 0060 DO 400 KL=1,LEV
15N 0061 DO 350 K=1,LASTLI
15N 0062 KK=LASTLI-K+1
15N 0063 IF(LINE(KK,KL).NE.1BLANK) GO TJ 375
15N 0064 350 CONTINUE
15N 0065 375 CONTINUE
15N 0066 IF(IFILEX.GT.0.AND.KL.EQ.LEV) WRITE(IFILE,1050) (LINE(I,KL),
15N 0067 I=1,KK)
15N 0069 IF(IFILEX.LT.0.AND.KL.EQ.1) WRITE(IFILE,1050) (LINE(I,KL),
15N 0071 I=1,KK)
15N 0071 IF(IFILEX.GT.0.AND.KL.NE.LEV) WRITE(IFILE,1060) (LINE(I,KL),
15N 0073 I=1,KK)
15N 0073 IF(IFILEX.LT.0.AND.KL.NE.1) WRITE(IFILE,1060) (LINE(I,KL),
15N 0075 I=1,KK)
15N 0075 400 CONTINUE

C
C WRITE LAST LINE (BORDER)
C
15N 0076 WRITE(IFILE,1070) (IBORDR,1=1,LASTLI)

C
C PRINT TRAILING SCALE INFORMATION
C
15N 0077 IF(NCPW.EQ.6) WRITE(IFILE,1030) (TITLE(I),I=1,NWORDS)
15N 0079 IF(NCPW.EQ.4) WRITE(IFILE,1001) (TITLE(I),I=1,NWORDS)
15N 0081 IF(NCPW.EQ.10) WRITE(IFILE,1002) (TITLE(I),I=1,NWORDS)
15N 0083 IF(NCPW.EQ.5) WRITE(IFILE,1033) (TITLE(I),I=1,NWORDS)
15N 0085 WRITE(IFILE,1010)
15N 0086 DELTA=(AMX-ANN)/FLOAT(NLEVEL)
15N 0087 DO 200 I=1,NLEVEL
15N 0088 XMIN=FLJAT(I-1)/FLOAT(NLEVEL)*(A4K-ANN)+ANN
15N 0089 XMAX=XMIN+DELTA
15N 0090 LEV=LEVEL(I)
15N 0091 DO 175 J=1,LEV
15N 0092 IF(IFILEX.GT.0.AND.J.EQ.LEV) WRITE(IFILE,1020) (CHARS(I,J),XMIN,

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LEVEL 2.3. (JUNE 78)      GREYSC      05/360      JETRAM H EXTENDED      DATE 80.199  1.57.13

ISN 0094      IF(1,1,EX,LT,0,AND,J,EQ,1) WRITE(1=ILE,1020) ICCHARS(1,J),KMIN,
ISN 0096      1 KMAX
ISN 0098      IF(1,1,FILEX,GT,0,AND,J,NE,LEV) WRITE(1=FILE,1030) ICCHARS(1,J)
ISN 0100      IF(1,1,FILEX,LT,0,AND,J,NE,1) WRITE(1=FILE,1030) ICCHARS(1,J)
ISN 0101      CONTINUE
ISN 0102      CONTINUE
ISN 0103      RETURN
ISN 0104      FORMAT(1=10,20A6)
ISN 0105      FORMAT(1=10,20A4)
ISN 0106      FORMAT(1=10,12A10)
ISN 0107      FORMAT(1=10,2AAS)
ISN 0108      FORMAT(1=10,32HGREY-SCALE CHARACTERS AND RANGES/1X)
ISN 0109      FORMAT(1=10,5X,E15.6,5X,E15.6)
ISN 0110      FORMAT(1=10,4X,A1)
ISN 0111      FORMAT(1=10,1X,20A6)
ISN 0112      FORMAT(1=10,1X,20A4)
ISN 0113      FORMAT(1=10,1X,12A10)
ISN 0114      FORMAT(1=10,1X,2AAS)
ISN 0115      FORMAT(1=10,132A1)
ISN 0116      FORMAT(1=10,132A1)
ISN 0117      FORMAT(1=10,132A1)
ISN 0118      FORMAT(1=10,132A1)
ISN 0119      END

*OPTIONS IN EFFECT*NAME(MAIN) OPTIMIZE(2) LINECOUNT(55) SIZE(MAX) AUTODRL(NONE)
*OPTIONS IN EFFECT*SOURCE EBCDIC NULIST MODECK OBJECT NUNAP N3FORMAT GOSTMT NOXRE= ALC NOANSF NOTERM IRM F1A
*STATISTICS*      SOURCE STATEMENTS =      118, PROGRAM SIZE =      4840, SUBPROGRAM NAME =GREYSC
*STATISTICS*      NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
80K BYTES OF CORE NOT USED

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LEVEL 2.3.. (JUNE 78)                                05/360      JATTAN H EXTENDED      DATE 80.199  1.57.15
REQUESTED OPTIONS:
OPTIONS IN EFFECT:  NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTOOBL(NONE)
SOURCE EBCDIC NOLIST NODECK OBJECT VJMAP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IDW FLA

ISM 0002      SUBROUTINE PRINTER(ISTEP)
ISM 0003      COMMON /CM2/REFNDX(16384),SY(128),CS(128),ZTS2(128),POSQ(128)
C              + .AMPARY(16384),RADARY(16384)
C              LEVEL 2:REFNDX,SN,CS,ZTSQ,PJSQ,AMPARY,RADARY
C              COMMON /ARRAYS/V(32768)
C              COMMON /PARAM/DZINC,MESH,LAMBOA,XD,FR,NQ,PCORP,OUTRAD,DX,NSTEPS,
C              + NOZIVC,MESHQ,MESHQ2,PI,WAVENM,XJ31,MS,MF,NF,MSHPTS
C              COMMON /PRINTPL/PGREY,PWAIST,PLTOST,PLTMAX,PLTOUT,IGREY
C              + .PLTFLD,PLTFLC
C              DIMENSION X(130),XWAIST(100),YWAIST(100),VRAD(100),WSTID(2),
C              + MAXI(2),MAXX(2),MAXY(2),WSFX(2),WSFY(2),XT(100)
C              DIMENSION FIELDS(130),XID(2),FLDID(2),FLOTTL(2)
C              EQUIVALENCE (MAXX,WSFX)
C              LOGICAL PGREY,PWAIST,PLTOST,PLTMAX,PLTFLD,PLTFLC,LAST
C              REAL -AM3DA
C              DATA NSTID/19H2ND MOMENTS VS DIST/
C              DATA VAXID/18H2PEAK IRRAD VS DIST/
C              DATA MAXX/10HAXIAL DIST/
C              DATA MAXY/15HPEAK IRRADIANCE/
C              DATA WSTY/17HBEAM WAIST RATIOS/
C              DATA XID/18HAXIAL DIST (MICR)/
C              DATA FLDID/10HIRRADIANCE/
C              DATA FLOTTL/20HIRRAD VS RADIAL DIST/
C              KSTEP=ISTEP+1
C              IF (PGREY.OR.PWAIST) WRITE(IOUT,2300) ISTEP
C              IF (PGREY.OR.PWAIST) PLTMAX=0,PLTOST=1 CALL PEAK(VMAX)
C              IF (.NOT.PWAIST.OR.PLTOST) GO TO 10
C              CALL SIZE(RADARY,MESH,MESH,XD,VJ,X2,Y2)
C              X2=1.41421356X2+DX
C              Y2=1.41421356Y2+DY
C              IF (PWAIST) WRITE(IOUT,2010) X2,Y2
C              IF (PWAIST) THIS IS STEP (12,/)
C              2000 FORMAT(17,14H THIS IS STEP (12,/)
C              2010 =OMAT(17,38H BEAM WAIST IN X AND Y DIRECTIONS IS (2(F10.5,2X),
C              + 7MICRONS)
C              10 CONTINUE
C              IF (PGREY) CALL GREYSC(IGREY,10,RADARY,MESH,MESH,MS,MF,1,NS,NF,1,
C              + 0,0,0,10-IRRADIANCE,10)
C              RETURN
C              END

ISM 0009
ISM 0010
ISM 0012
ISM 0014
ISM 0016
ISM 0018
ISM 0019
ISM 0020
ISM 0021
ISM 0022
ISM 0023
ISM 0024
ISM 0026
ISM 0027

OPTIONS IN EFFECT: NAME(MAIN) OPTIMIZE(2) LINECOUNT(56) SIZE(MAX) AUTOOBL(NONE)
OPTIONS IN EFFECT: SOURCE EBCDIC NOLIST NODECK OBJECT VJMAP NOFORMAT GOSTMT NOXREF ALC NOANSF NOTERM IDW FLA
STATISTICS* SOURCE STATEMENTS = 26. PROGRAM SIZE = 692. SUBPROGRAM NAME =PRNTEP
STATISTICS* NO DIAGNOSTICS GENERATED
***** END OF COMPILATION *****
112K BYTES OF CORE NOT USED

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DATE 80.19.21.57.15

OS/360 F33TRAN M EXTENDED

LEVEL 2.3.0 (JUNE 78)

STATISTICS NO DIAGNOSTICS THIS STEP

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IEF1421 - S..P EXECUTED - COND CODE 0000
IEF2851 SYS1.FORTRAN PASSED
IEF2851 VOL SER NOS= ACS002.
IEF2851 SYS80199.T215654.RV000.SBI404F2.LOADSET PASSED
IEF2851 VOL SER NOS= SYSDA1.
IEF2851 SYS80199.T215654.RV000.SBI404F2.R0000001 DELETED
IEF2851 VOL SER NOS= SYSDA3.
IEF2851 SYS80199.T215654.RV000.SBI404F2.R0000002 DELETED
IEF2851 VOL SER NOS= SYSDA2.
IEF2851 SYS80199.T215654.SV000.SBI404F2.R0000003 SYSQUT
IEF2851 VOL SER NOS= SYSDA3.
IEF2851 SYS80199.T215654.RV000.SBI404F2.S0000004 SYSIN
IEF2851 VOL SER NOS= SYSDA2.
IEF2851 SYS80199.T215654.RV000.SBI404F2.S0000004 DELETED
IEF2851 VOL SER NOS= SYSDA2.

CCN3011 SYSRINT PRINTED 24 TOTAL PAGES. REQUIRED 10 TRACKS.

REGION ALLOCATION CORE USED
250K

STEP CPU TIME JOB CPU TIME STEP I/O CJVT JOB I/O COUNT
3.14S 3.14S 306 306

XXGO EXEC PGM=GLoader,COND=(5,LT,FORT),TIME=2TG,REGION=2RG 00000320
IEF6531 SUBSTITUTION JCL - PGM=LLADER,COND=(5,LT,FORT),TIME=1439,REGION=450K
XXFT05F001 DD DDNAME=SYSIN 00000340
XXSTEPL1B DD DISP=(SHR,PASS),DSN=SYSLEVEL,*,FORTRAN 00000360
IEF6531 SUBSTITUTION JCL - DISP=(SHR,PASS),DSN=SYS1.FORTAN 00000390
XXSYSL1B DD DISP=(SHR,PASS),DSN=SYSLEVEL,*,FORTRAN 00000400
IEF6531 SUBSTITUTION JCL - DISP=(SHR,PASS),DSN=SYS1.FORT-13 00000420
XXSYSLIN DD DSN=GLLOADSET,DISP=(OLD,DELETE) 00000440
XXSYSPROUT DD SYSDA=A,SPACE=(133,(200,75),RLSE) 00000460
XXFT06F001 DD SYSDA=A,UNIT=SYSDA,SPACE=(TRK,6PG,RLSE), 00000480
IEF6531 SUBSTITUTION JCL - SYSDA=A,UNIT=SYSDA,SPACE=(TRK,25,RLSE), 00000500
XX DCB=(RECFM=FBA,--RECL=133,BLKSIZE=4123)
//GO.SYSIN DD *
//
IEF2361 ALLOC. FOR SBI404F2 GD
IEF2371 153 ALLOCATED TO FT05F001
IEF2371 822 ALLOCATED TO STEPL1B
IEF2371 822 ALLOCATED TO SYS-1B
IEF2371 153 ALLOCATED TO SYSLIN
IEF2371 713 ALLOCATED TO SYSPROUT
IEF2371 265 ALLOCATED TO FT05F001

```


OPTIONS USED - PRINT,MAP,NOLIST,CALL,RES,NOTERM,SIZE=3223,NAME=00GD

[illegible]

P. 10. 17710x

[illegible]

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REFNDX

GREY-SCALE CHARACTERS AND RANGES

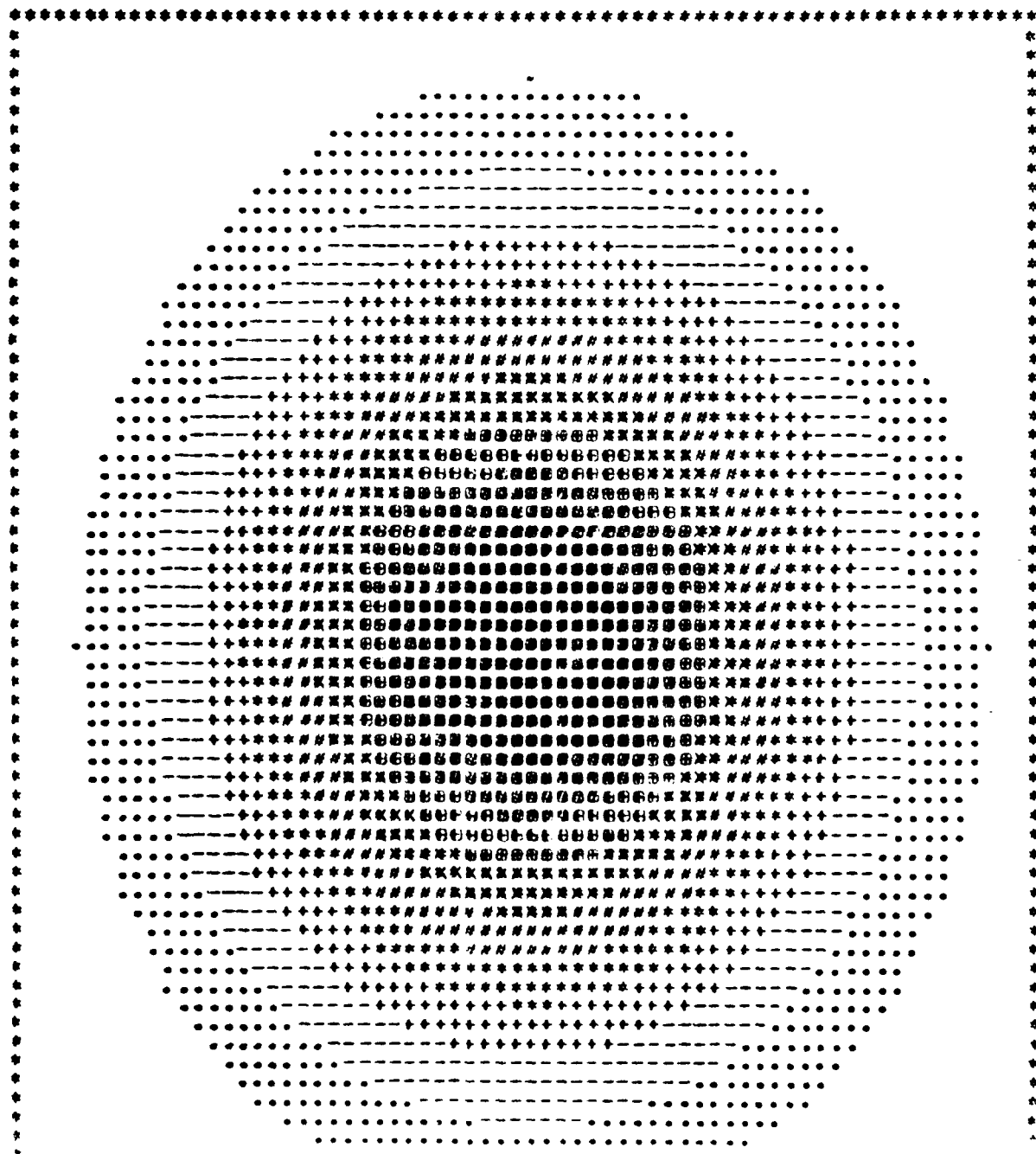
.	0.145928E+01	0.146135E+01
-	0.146135E+01	0.146342E+01
+	0.146342E+01	0.146549E+01
*	0.146550E+01	0.146757E+01
#	0.146757E+01	0.146964E+01
/	0.146964E+01	0.147171E+01
\	0.147171E+01	0.147378E+01
0	0.147378E+01	0.147585E+01
1	0.147585E+01	0.147793E+01
2	0.147793E+01	0.148000E+01

THIS IS STEP 0

SUML = 0.4860770E+00 SU4R = 0.5139230E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.99481 19.99631 MICRONS

IRRADIAN



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IRRADIANCE

GREY-SCALE CHARACTERS AND RANGES

.	0.597604E-02	0.105378E+00
-	0.105378E+00	0.204781E+00
+	0.204781E+00	0.304183E+00
*	0.304183E+00	0.403585E+00
#	0.403585E+00	0.502988E+00
%	0.502988E+00	0.602390E+00
@	0.602390E+00	0.701793E+00
!	0.701793E+00	0.801195E+00
"	0.801195E+00	0.900598E+00
#	0.900598E+00	0.100000E+01

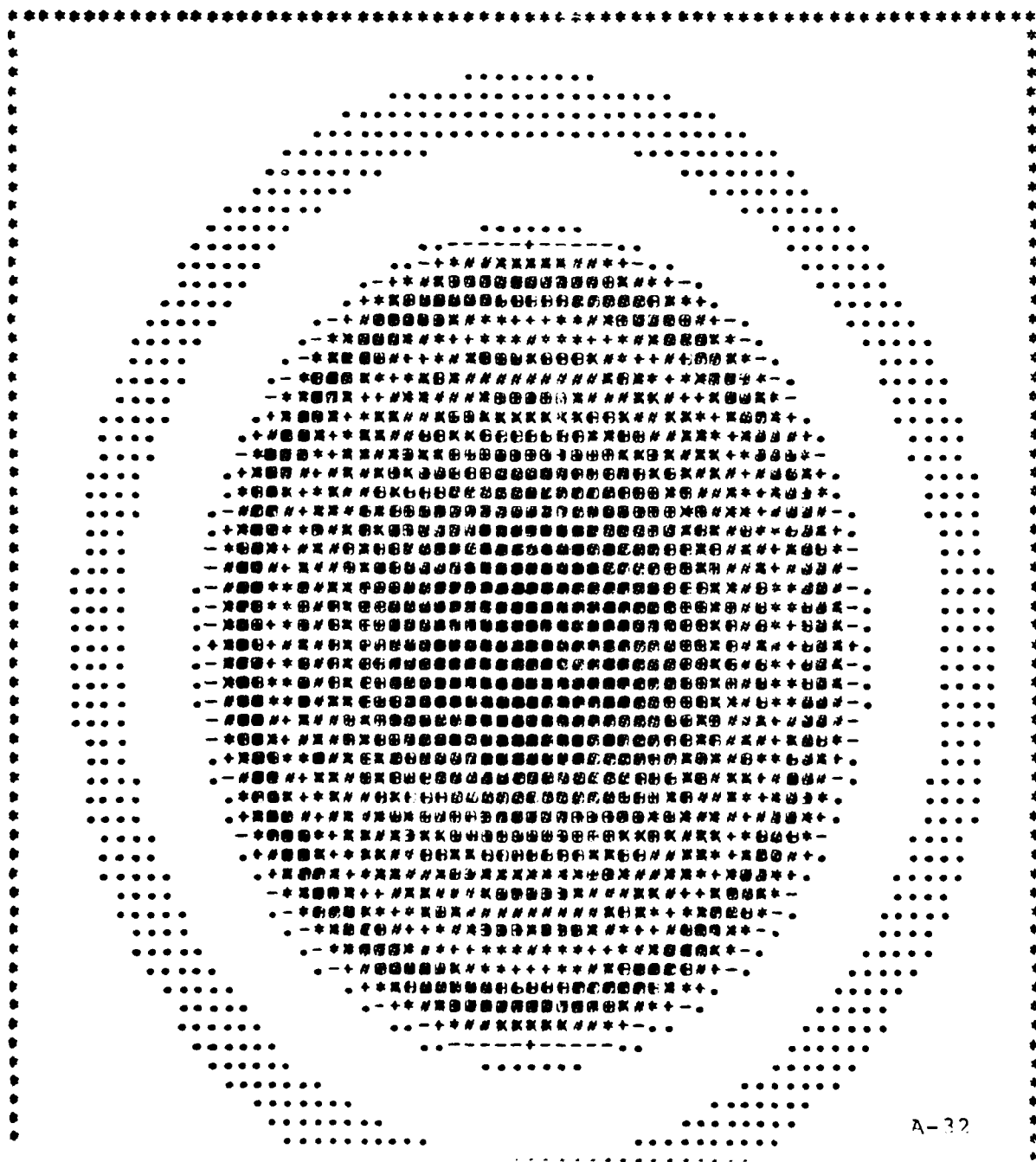
THIS IS STEP 1

SUML = 0.4859293E+00 SUMR = 0.5140707E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.50356 19.50494 MICRONS

IRRADIATION

P. 168



A-32

ERRAD (AN

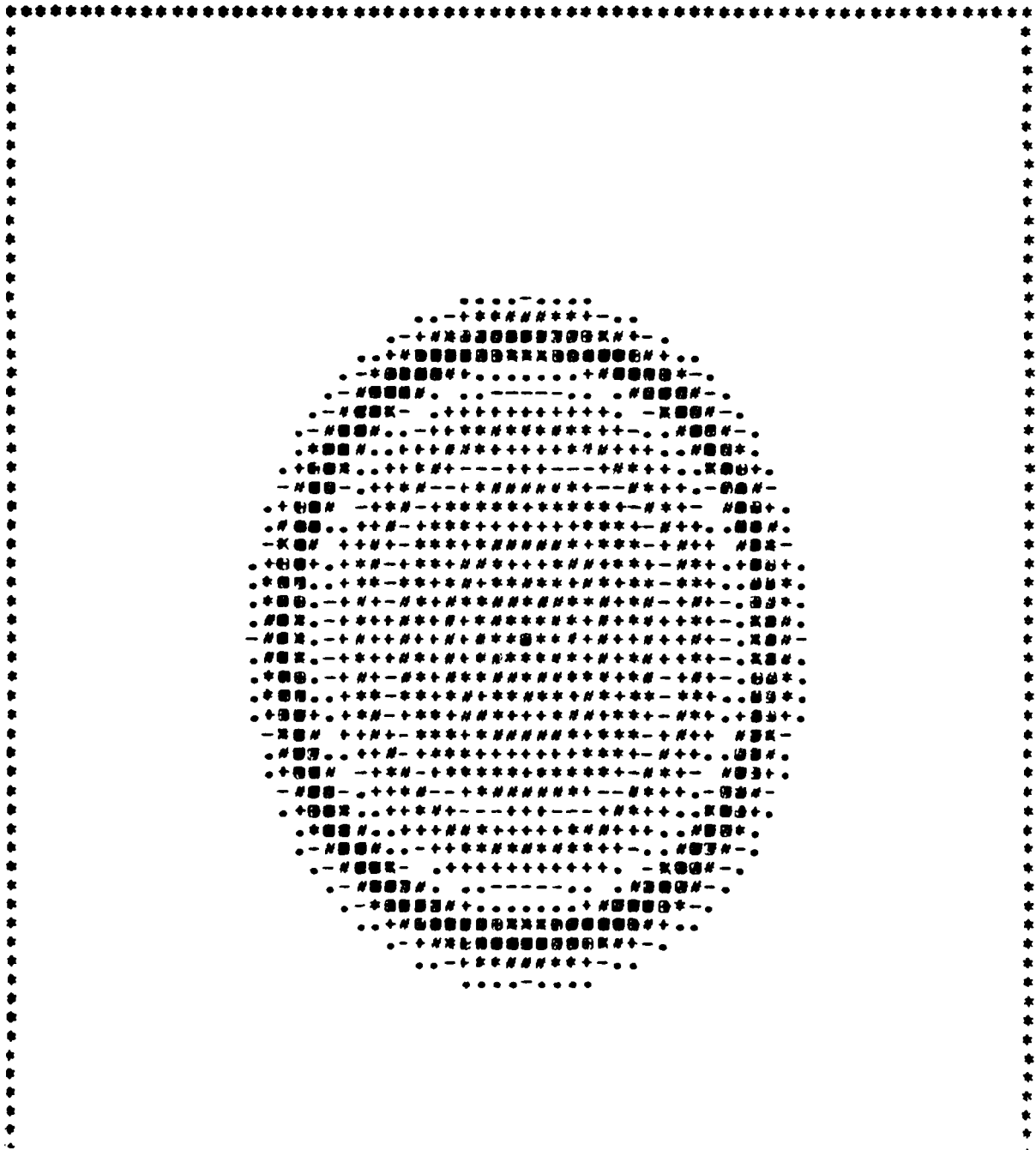
0.585011E-02	0.105238E+00
0.105238E+00	0.204626E+00
0.204626E+00	0.304013E+00
0.304013E+00	0.403401E+00
0.403401E+00	0.502789E+00
0.502789E+00	0.602177E+00
0.602177E+00	0.701565E+00
0.701565E+00	0.800952E+00
0.800952E+00	0.900340E+00
0.900340E+00	0.999728E+00

SUML = 0.4851605E+00 SU42 = 0.5148394E+00

Λ-33

P-100

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.323949E-02	0.211144E+00
-	0.211144E+00	0.419048E+00
+	0.419048E+00	0.626953E+00
*	0.626953E+00	0.834857E+00
#	0.834857E+00	0.104276E+01
%	0.104276E+01	0.125066E+01
^	0.125067E+01	0.145857E+01
^	0.145857E+01	0.166647E+01
^	0.166647E+01	0.187438E+01
^	0.187438E+01	0.208228E+01

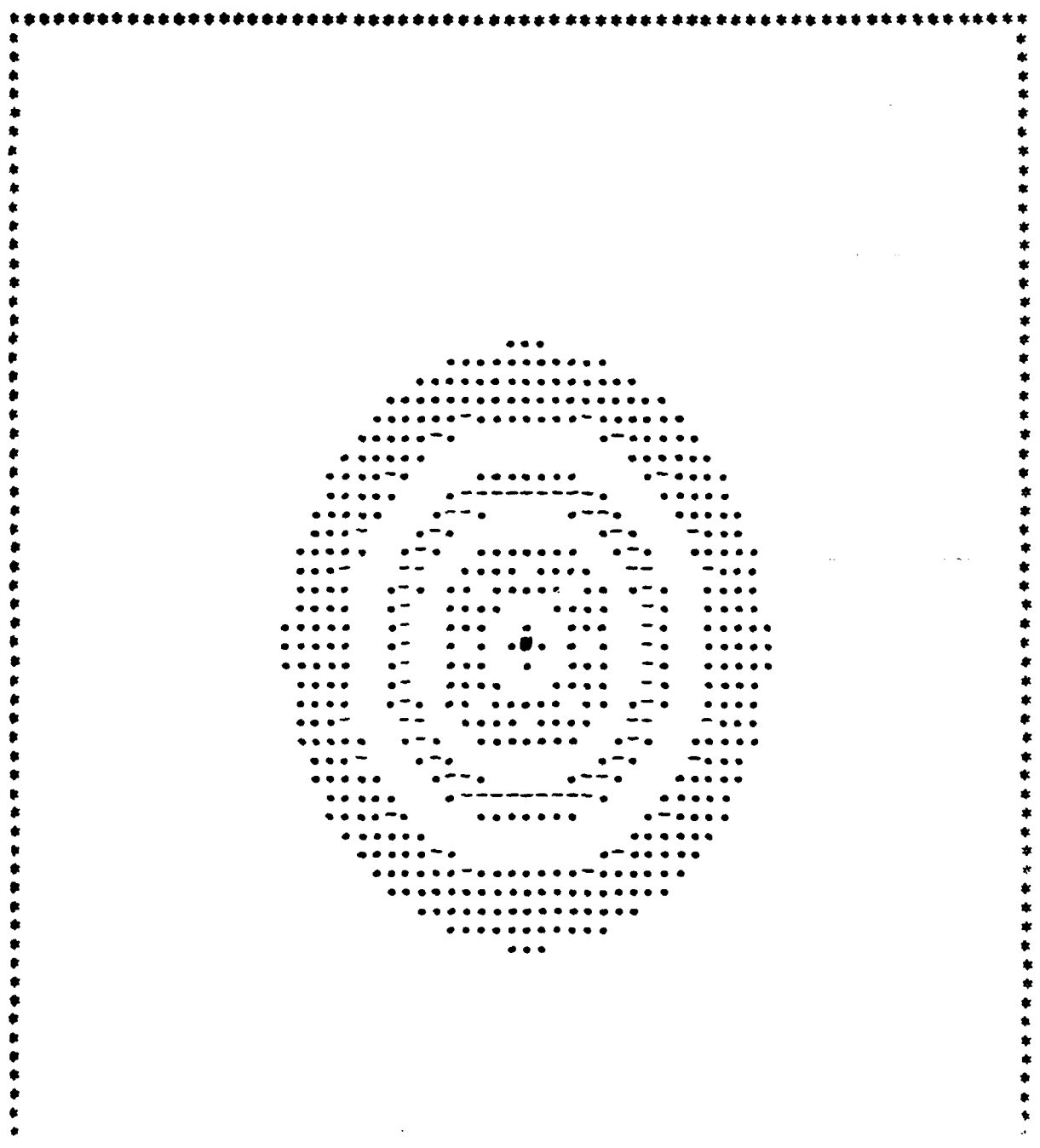
THIS IS STEP 3

SUNL = 0.4817133E+00 SUMR = 0.5182866E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.07376 18.07498 MICRONS

P-102

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.759699E-05	0.842067E+00
-	0.842067E+00	0.168413E+01
+	0.168412E+01	0.252618E+01
*	0.252618E+01	0.336824E+01
#	0.336824E+01	0.421030E+01
%	0.421030E+01	0.505236E+01
^	0.505236E+01	0.589442E+01
^	0.589442E+01	0.673648E+01
^	0.673648E+01	0.757854E+01
^	0.757854E+01	0.842060E+01

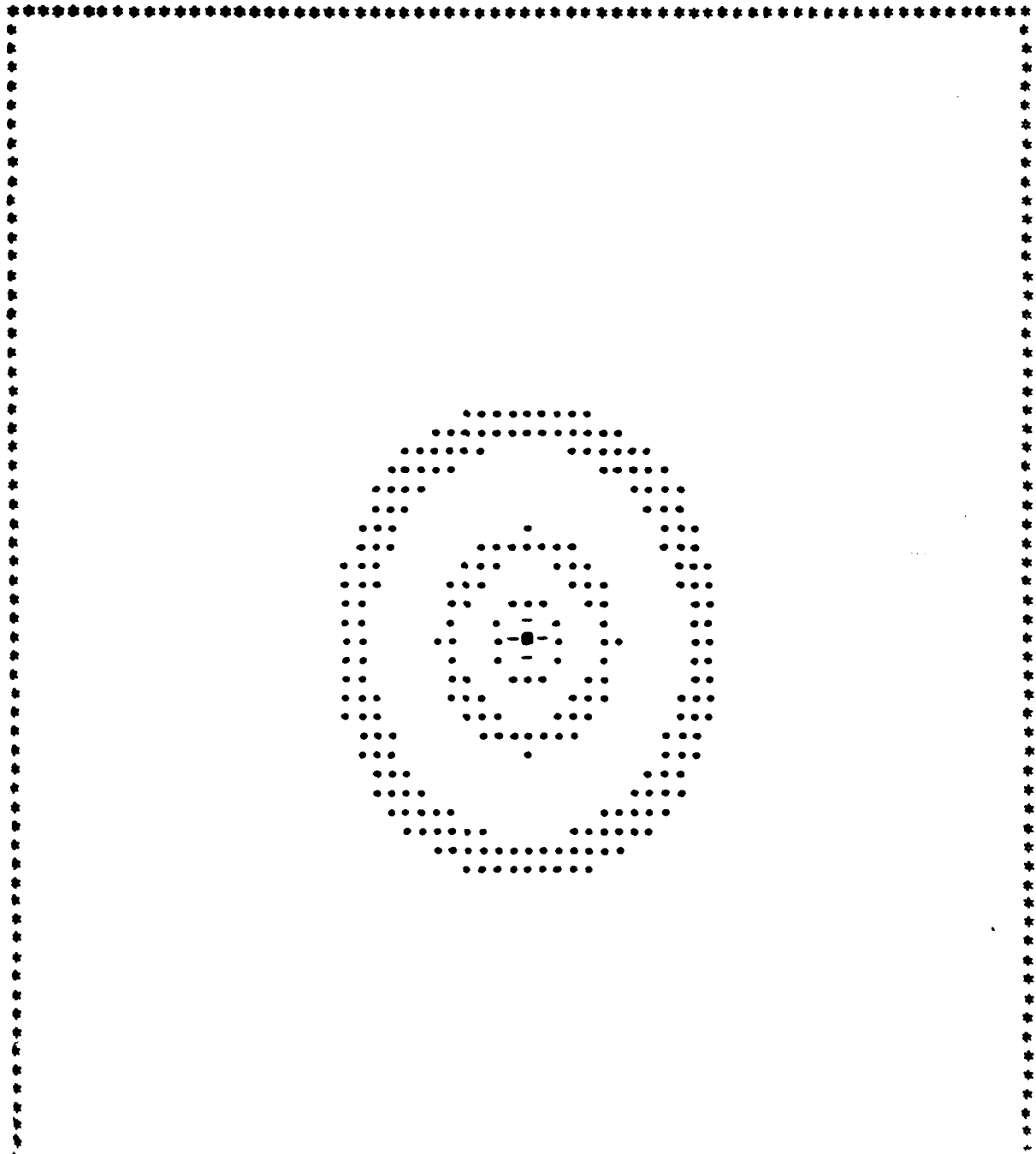
THIS IS STEP 4

SUML = 0.4754750E+00 SUMR = 0.5245249E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.64449 17.64572 MICRONS

P-10

IRRADIAN



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IRRADIANCE

GREY-SCALE CHARACTERS AND RANGES

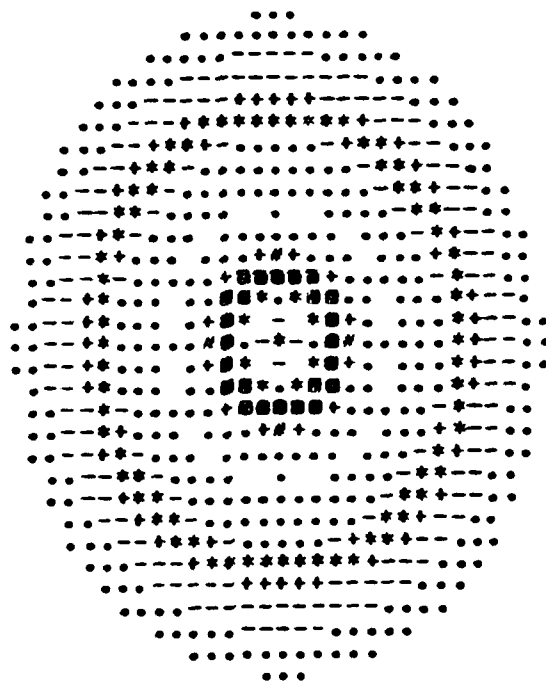
.	0.247009E-03	0.153054E+01
-	0.153054E+01	0.306083E+01
+	0.306083E+01	0.459112E+01
*	0.459112E+01	0.612142E+01
*	0.612142E+01	0.765171E+01
#	0.765171E+01	0.918200E+01
%	0.918200E+01	0.107123E+02
■	0.107123E+02	0.122426E+02
■	0.122426E+02	0.137729E+02
■	0.137729E+02	0.153032E+02

THIS IS STEP 5

SUML = 0.4803946E+00 SUMR = 0.5196053E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.43616 17.43765 MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.348578E-05	0.546719E+00
-	0.546719E+00	0.109343E+01
+	0.109343E+01	0.154015E+01
*	0.164015E+01	0.218686E+01
#	0.218696E+01	0.273358E+01
%	0.273358E+01	0.328029E+01
@	0.328029E+01	0.382701E+01
0	0.382701E+01	0.437372E+01
1	0.437373E+01	0.492044E+01
2	0.492044E+01	0.546716E+01

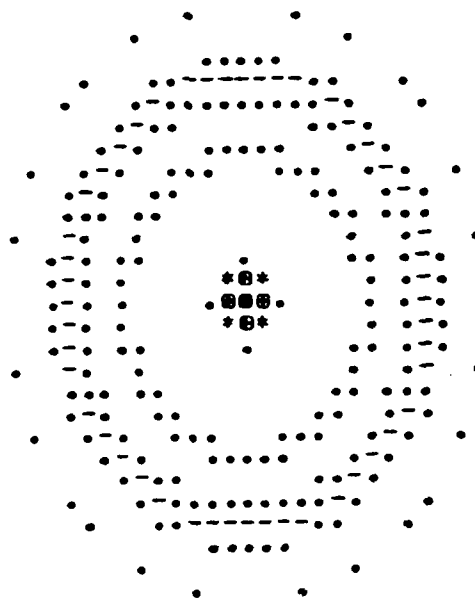
THIS IS STEP 6

SUML = 0.4704280E+00 SUMR = 0.5295719E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.47115 17.47259 MICRONS

P. 124

IRRADIAN



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ERRADIAN

GREY-SCALE CHARACTERS AND RANGES

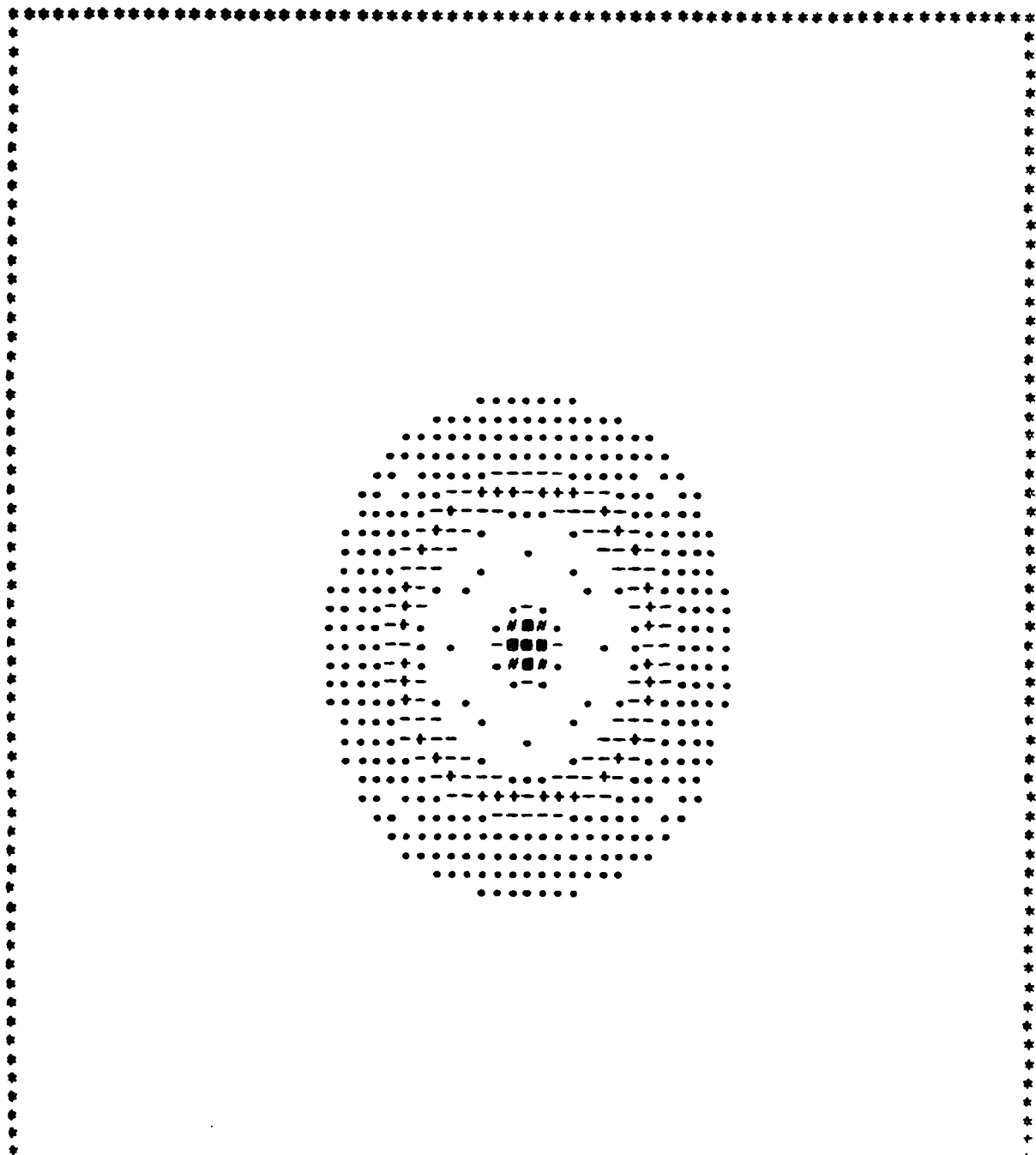
.	0.341598E-04	0.146743E+01
-	0.146743E+01	0.293483E+01
+	0.293483E+01	0.440223E+01
*	0.440223E+01	0.586963E+01
#	0.586963E+01	0.733703E+01
%	0.733703E+01	0.880443E+01
@	0.880443E+01	0.102718E+02
0	0.102718E+02	0.117392E+02
1	0.117392E+02	0.132066E+02
2	0.132066E+02	0.146740E+02

THIS IS STEP 7

SUML = 0.4736281E+00 SUMR = 0.5263718E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 17.75745 17.75934 MICRONS

IRRADIATION

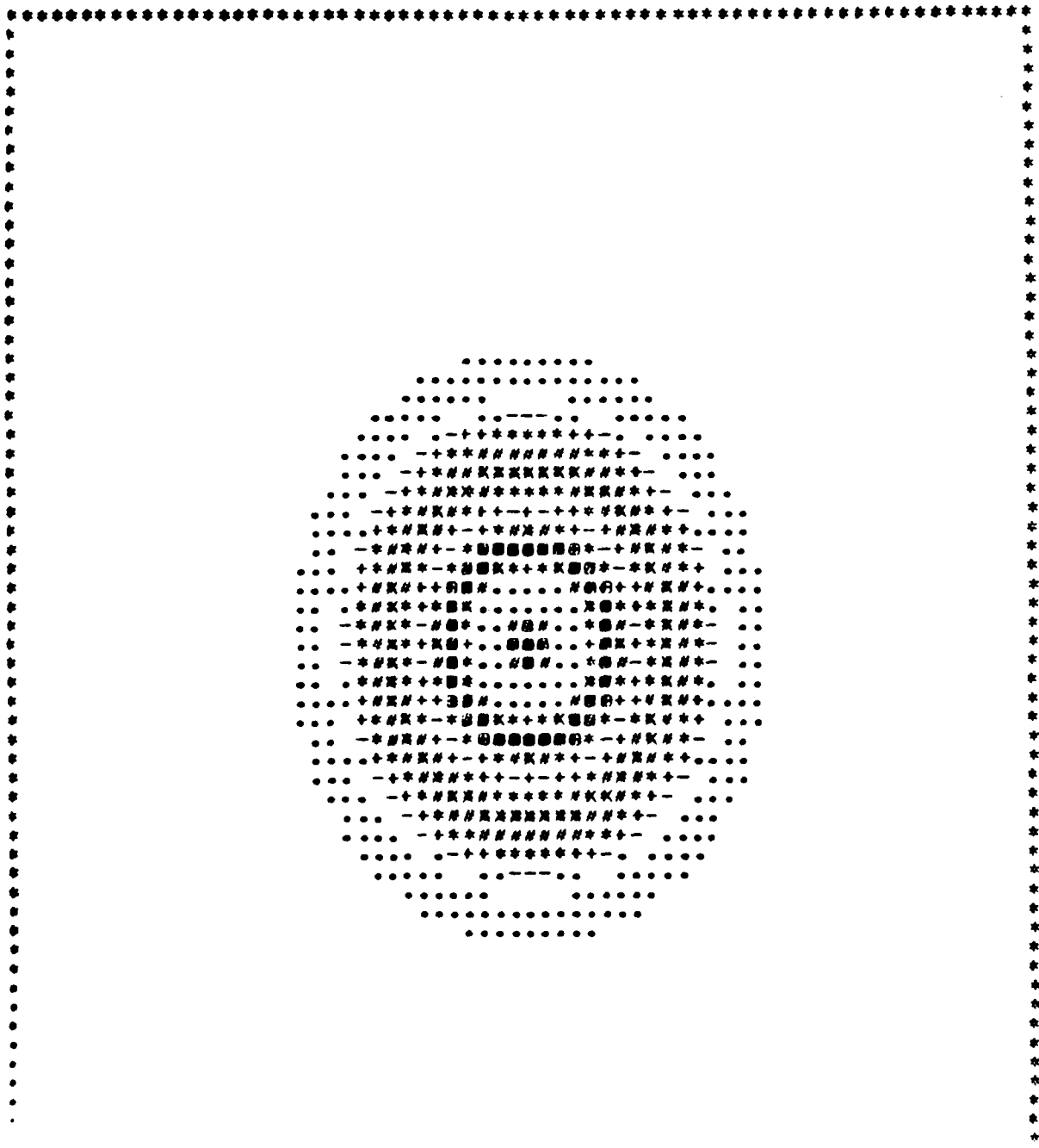


GREY-SCALE CHARACTERS AND RANGES

THIS IS STEP 8

BEAM WAIST IN X AND Y DIRECTIONS IS 18.20370 18.20607 MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

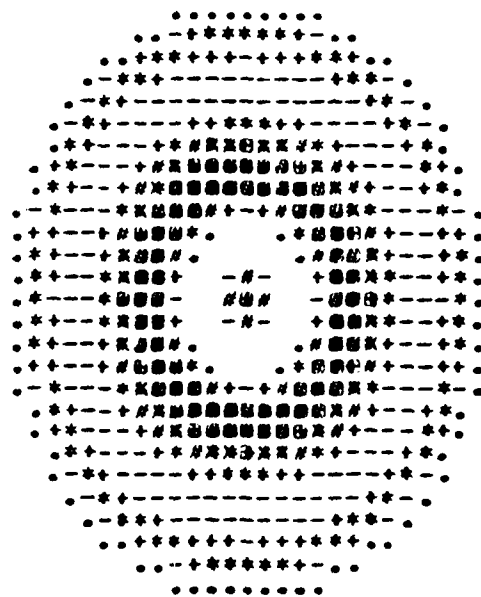
.	0.213262E-03	0.358114E+00
-	0.358113E+00	0.716014E+00
+	0.716014E+00	0.107391E+01
*	0.107391E+01	0.143181E+01
#	0.143181E+01	0.178971E+01
%	0.178971E+01	0.214761E+01
&	0.214761E+01	0.250551E+01
'	0.250551E+01	0.286341E+01
(0.286341E+01	0.322131E+01
)	0.322132E+01	0.357922E+01

THIS IS STEP 9

SUML = 0.4812048E+00 SUMR = 0.5187951E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.54669 19.54912 MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

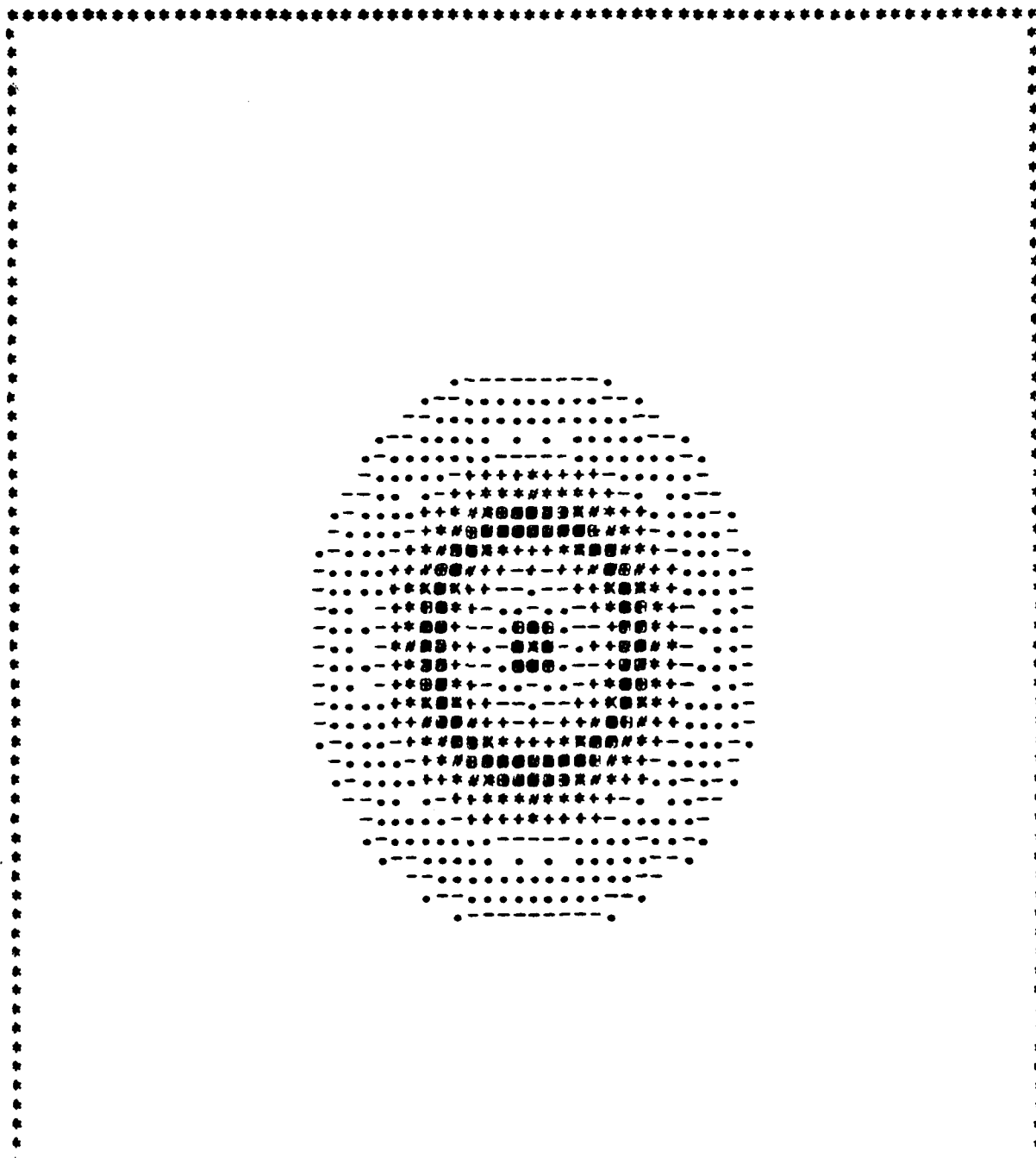
.	0.753595E-04	0.379830E+00
-	0.379830E+00	0.759584E+00
+	0.759584E+00	0.113934E+01
*	0.113934E+01	0.151909E+01
#	0.151909E+01	0.189885E+01
0	0.189885E+01	0.227860E+01
1	0.227860E+01	0.265836E+01
2	0.265836E+01	0.303811E+01
3	0.303811E+01	0.341786E+01
4	0.341786E+01	0.379762E+01

THIS IS STEP 10

SUM1 = 0.4789906E+00 SUM2 = 0.5210093E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.63527 18.63782 MICRONS

IRRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

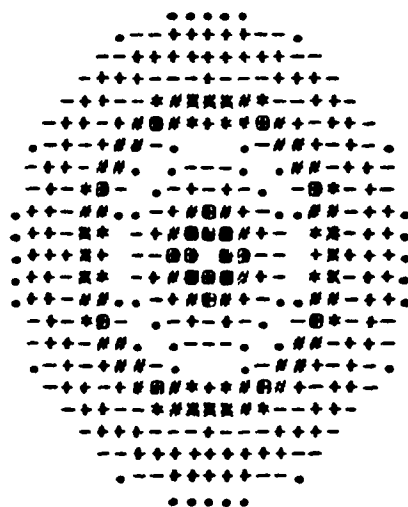
.	0.103213E-03	0.410842E+00
-	0.410842E+00	0.821581E+00
+	0.821582E+00	0.123232E+01
*	0.123232E+01	0.164306E+01
#	0.164306E+01	0.205380E+01
@	0.205380E+01	0.246454E+01
A	0.246454E+01	0.287528E+01
B	0.287528E+01	0.328602E+01
C	0.328602E+01	0.369676E+01
D	0.369676E+01	0.410749E+01

THIS IS STEP 11

SUML = 0.4775044E+00 SUM2 = 0.5224955E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.77403 18.77573 MICRONS

ERRADIAN



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ERRADIAN

GREY-SCALE CHARACTERS AND RANGES

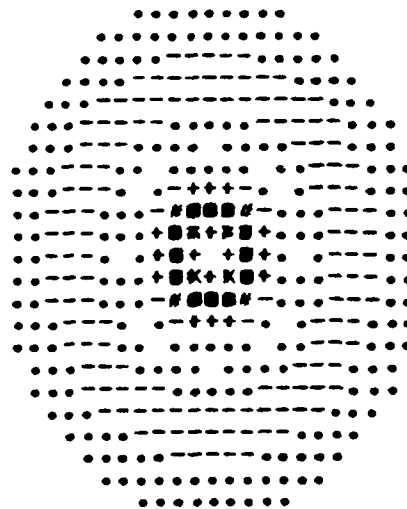
.	0.373785E-04	0.594157E+00
-	0.594157E+00	0.118828E+01
+	0.118828E+01	0.178239E+01
*	0.178239E+01	0.237651E+01
#	0.237651E+01	0.297063E+01
%	0.297063E+01	0.356475E+01
@	0.356475E+01	0.415887E+01
0	0.415887E+01	0.475299E+01
1	0.475299E+01	0.534711E+01
2	0.534711E+01	0.594123E+01

THIS IS STEP 12

SUML = 0.4750217E+00 SUMR = 0.5249783E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.92630 18.92923 MICRONS

IRRADIANT



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

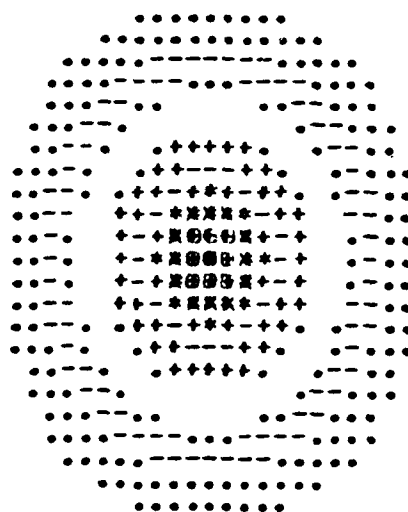
.	0.126287E-04	0.898516E+00
-	0.898516E+00	0.179702E+01
+	0.179702E+01	0.269552E+01
*	0.269552E+01	0.359403E+01
#	0.359403E+01	0.449253E+01
%	0.449253E+01	0.539103E+01
@	0.539103E+01	0.628954E+01
^	0.628954E+01	0.718804E+01
~	0.718804E+01	0.808654E+01
▯	0.808654E+01	0.898505E+01

THIS IS STEP 13

SUML = 0.4702227E+00 SUMR = 0.5297772E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.73090 19.73331 MICRONS

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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.703742E-04	0.848021E+00
-	0.848021E+00	0.169597E+01
+	0.169597E+01	0.254392E+01
*	0.254392E+01	0.339187E+01
#	0.339187E+01	0.423982E+01
%	0.423982E+01	0.508778E+01
^	0.508778E+01	0.593573E+01
&	0.593573E+01	0.678368E+01
@	0.678368E+01	0.763163E+01
0	0.763163E+01	0.847958E+01

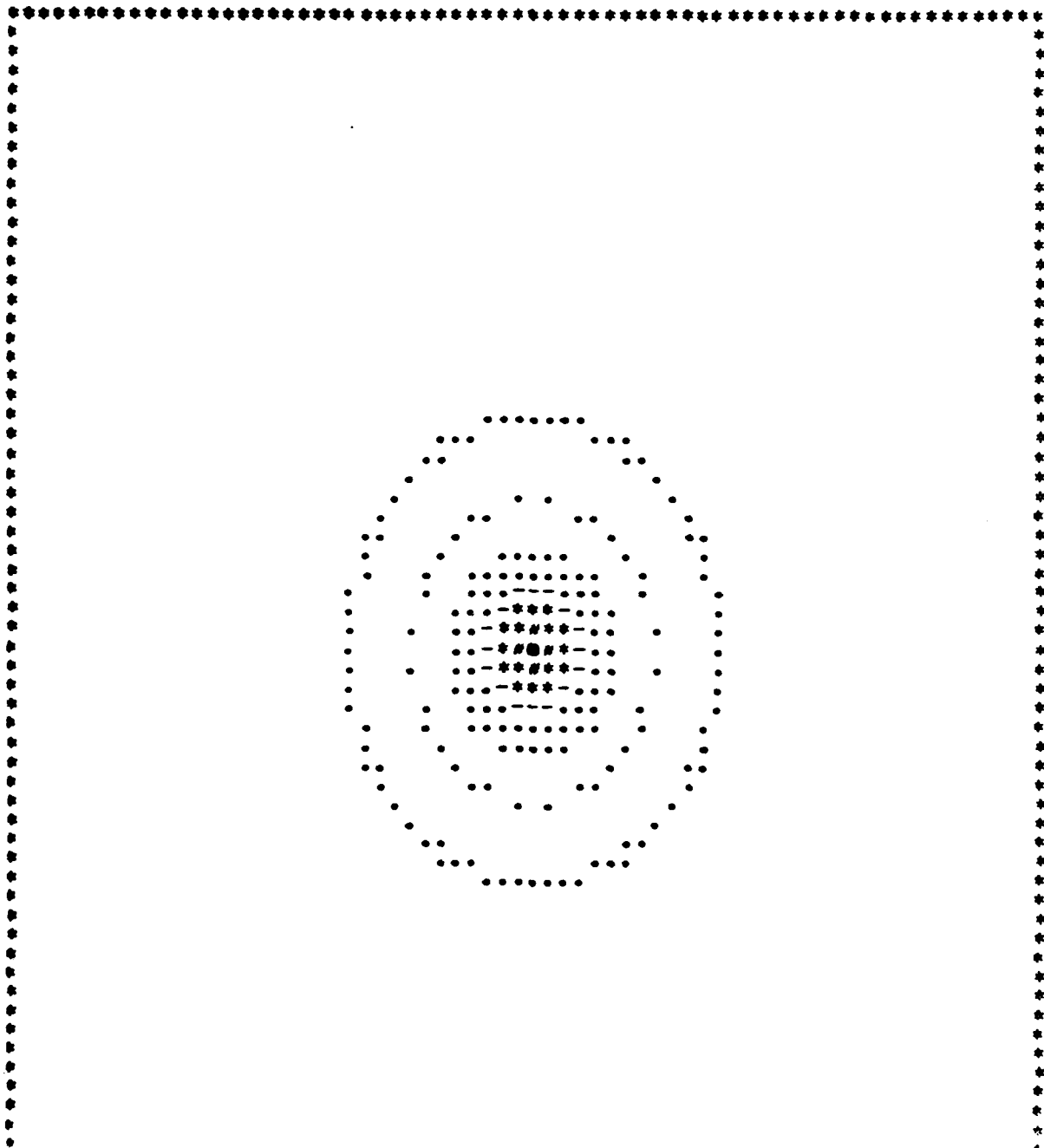
THIS IS STEP 14

SUML = 0.4651438E+00 SUMR = 0.5348561E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.77950 18.78253 MICRONS

2100

ERRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.244114E-03	0.155652E+01
-	0.155652E+01	0.311279E+01
+	0.311279E+01	0.466906E+01
*	0.466906E+01	0.622533E+01
#	0.622534E+01	0.778161E+01
%	0.778161E+01	0.933788E+01
0	0.933788E+01	0.108942E+02
1	0.108942E+02	0.124504E+02
2	0.124504E+02	0.140067E+02
3	0.140067E+02	0.155630E+02

THIS IS STEP 15

SUML = 0.4690364E+00 SUMR = 0.5309635E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 18.87981 18.88269 MICRONS

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EMTEC ENGINEERING INC LOS ANGELES CA
THEORETICAL ANALYSIS OF MULTIMODE FIBER STRUCTURES.(U)
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UNCLASSIFIED

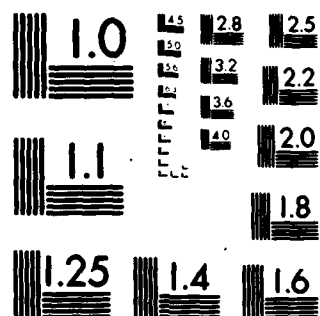
RADC-TR-80-332

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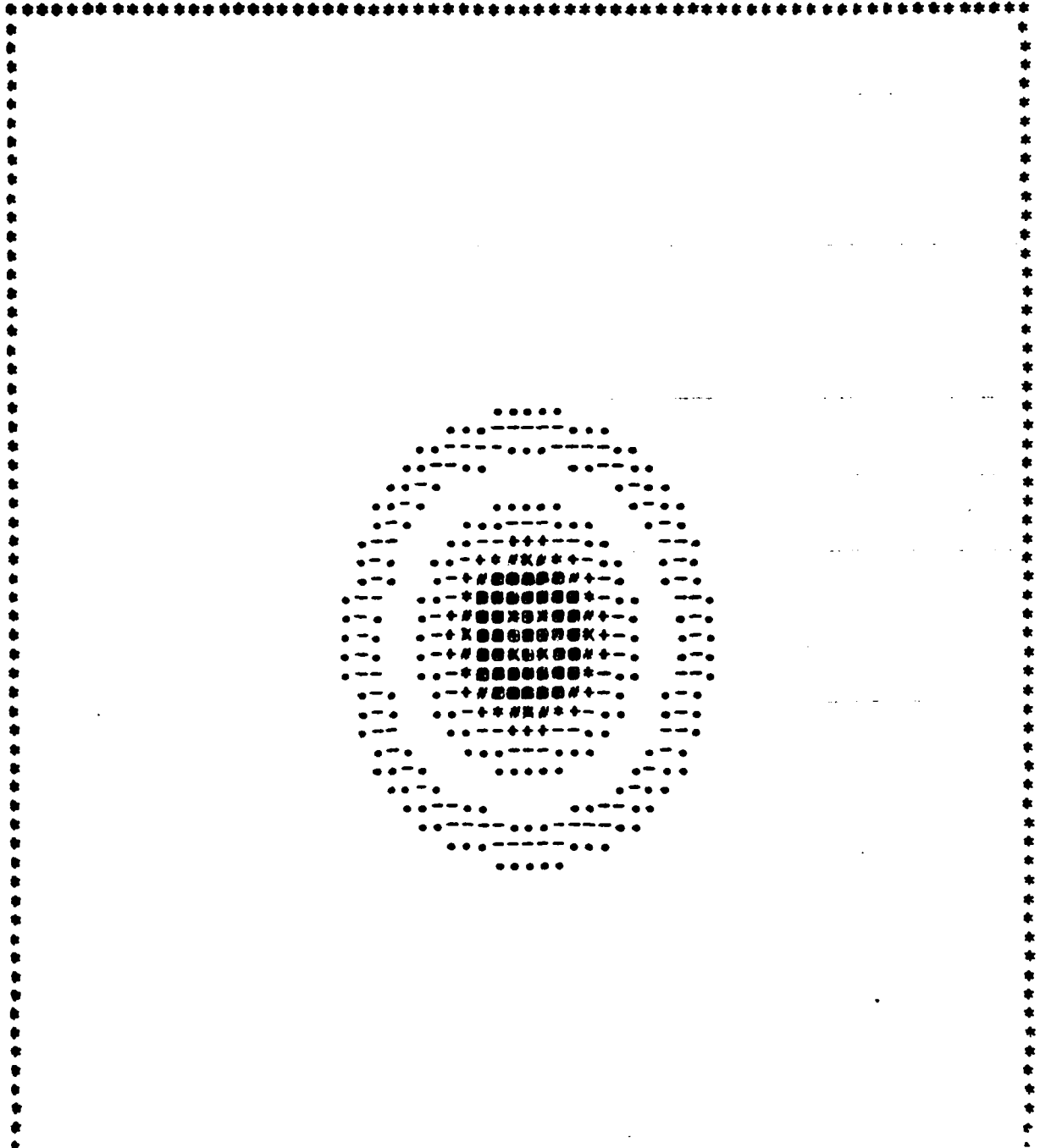
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IRRADIAN



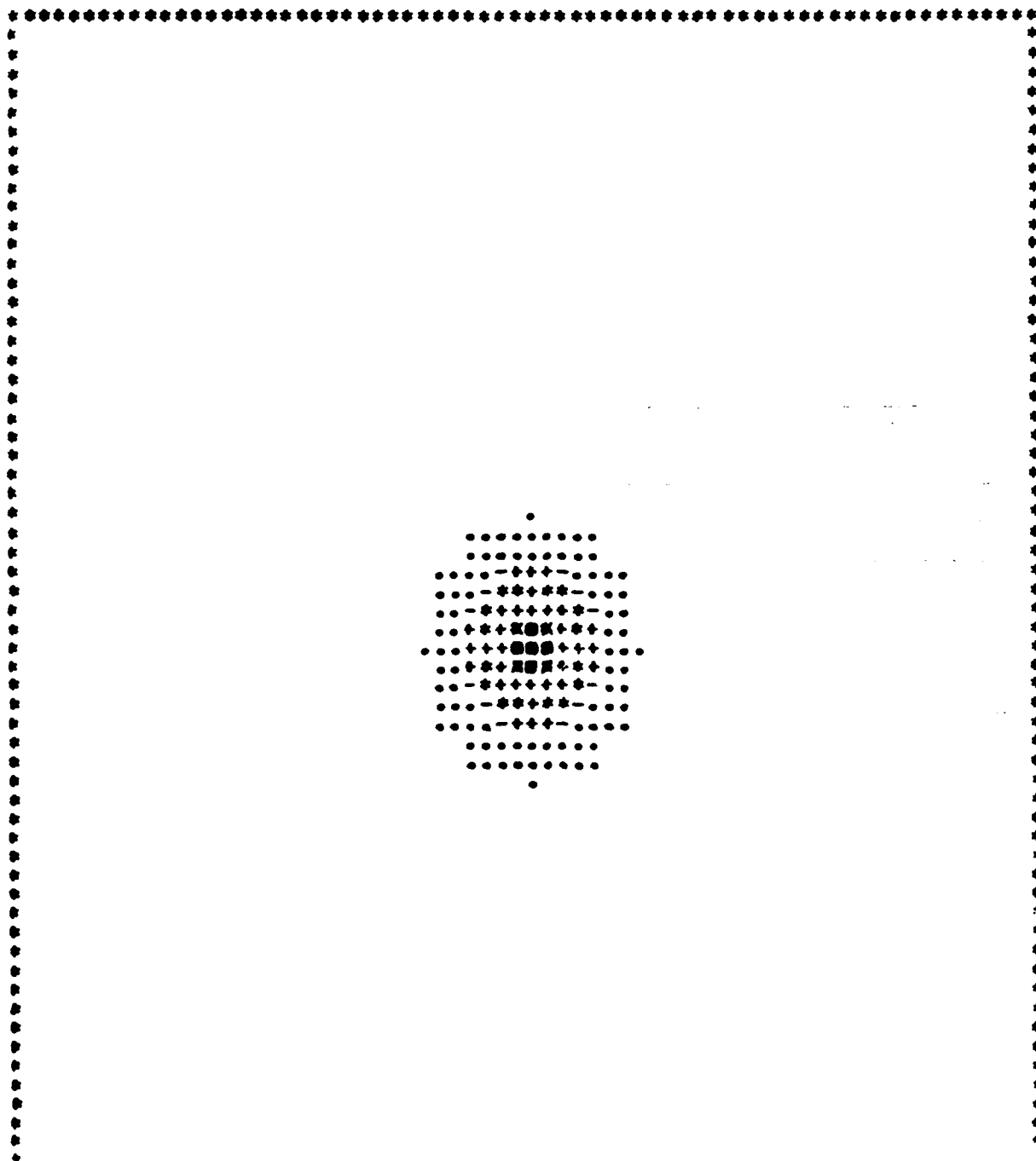
GREY-SCALE CHARACTERS AND RANGES

	0.246905E-03	0.697666E+00
•	0.697666E+00	0.139509E+01
-	0.139508E+01	0.209250E+01
+	0.209250E+01	0.278992E+01
*	0.278992E+01	0.348734E+01
#	0.348734E+01	0.418476E+01
\$	0.418476E+01	0.488218E+01
%	0.488218E+01	0.557960E+01
^	0.557960E+01	0.627702E+01
&	0.627702E+01	0.697444E+01

SUML = 0.4632549E+00 SUMR = 0.5367450E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.10667 19.10956 MICRONS

IRRADIAN



ERRADIAN

GREY-SCALE CHARACTERS AND RANGES

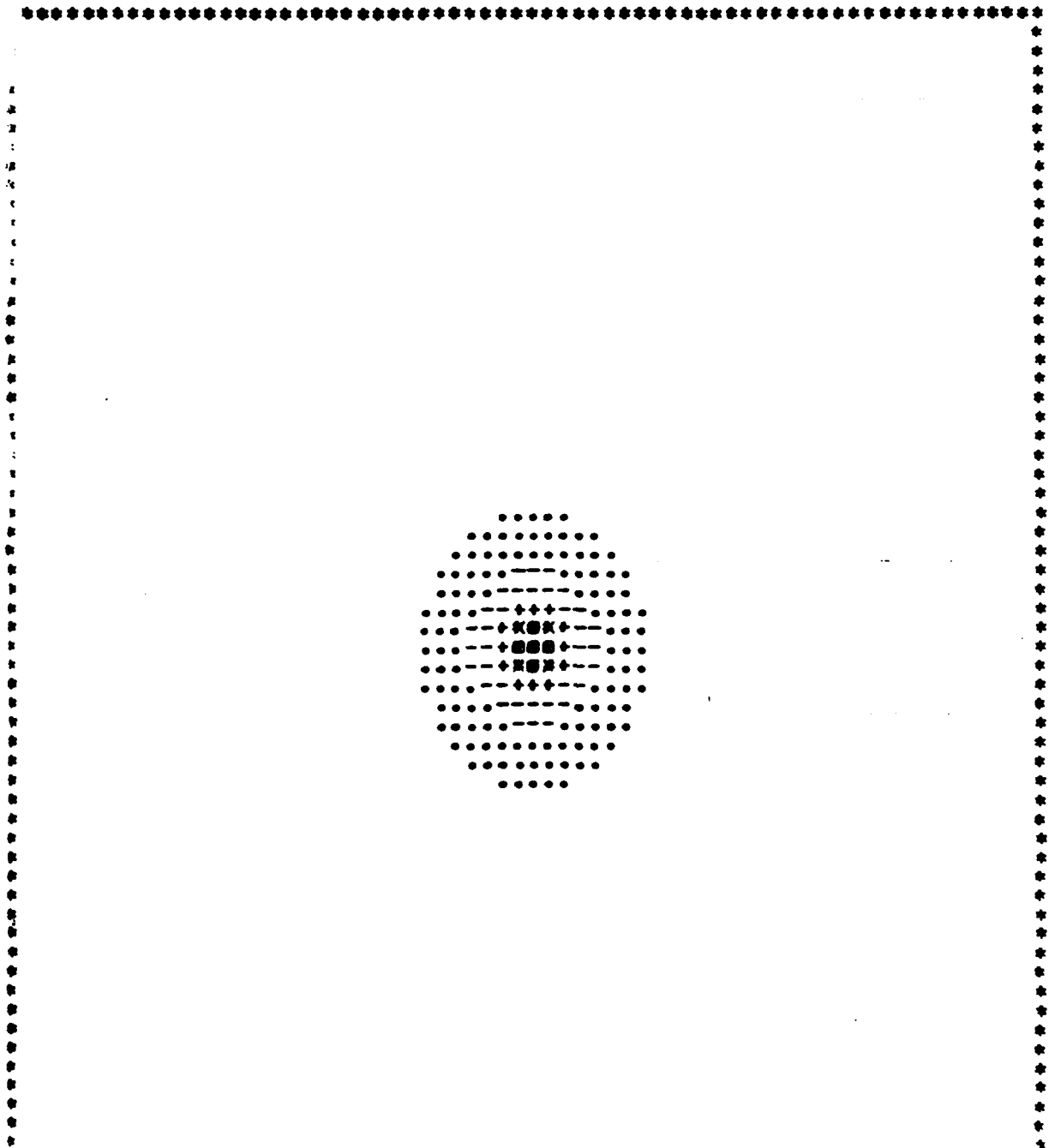
.	0.911394E-03	0.131676E+01
-	0.131676E+01	0.263261E+01
+	0.263261E+01	0.394845E+01
*	0.394845E+01	0.526430E+01
#	0.526430E+01	0.658015E+01
%	0.658015E+01	0.789600E+01
@	0.789600E+01	0.921185E+01
!	0.921185E+01	0.105277E+02
"	0.105277E+02	0.118435E+02
#	0.118435E+02	0.131594E+02

THIS IS STEP 17

SUML = 0.4618073E+00 SU42 = 0.5381926E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.16985 19.17290 MICRONS

ERRADIAN



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IRRADIAN

GREY-SCALE CHARACTERS AND RANGES

.	0.192583E-04	0.159542E+01
-	0.159542E+01	0.319081E+01
+	0.319081E+01	0.478621E+01
*	0.478621E+01	0.638161E+01
#	0.638161E+01	0.797700E+01
%	0.797700E+01	0.957240E+01
^	0.957240E+01	0.111678E+02
~	0.111678E+02	0.127632E+02
0	0.127632E+02	0.143586E+02
1	0.143586E+02	0.159540E+02

THIS IS STEP 18

SUML = 0.4730968E+00 SUMR = 0.5269032E+00

BEAM WAIST IN X AND Y DIRECTIONS IS 19.34006 19.34309 MICRONS

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RADC plans and executes research, development, test and selected acquisition programs in support of Command, Control Communications and Intelligence (C³I) activities. Technical and engineering support within areas of technical competence is provided to ESD Program Offices (POs) and other ESD elements. The principal technical mission areas are communications, electromagnetic guidance and control, surveillance of ground and aerospace objects, intelligence data collection and handling, information system technology, ionospheric propagation, solid state sciences, microwave physics and electronic reliability, maintainability and compatibility.

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RADC-TR-68-100